

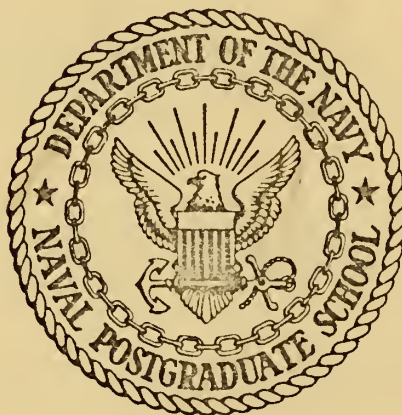
CHARACTERIZATION OF ELASTIC SOLIDS
USING
FINITE ELEMENT METHODS

Brian Alfred Edwards

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THESIS

CHARACTERIZATION OF ELASTIC SOLIDS
USING
FINITE ELEMENT METHODS

by

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Thesis Advisor:

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December 1972

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Characterization of Elastic Solids
Using
Finite Element Methods

by

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ABSTRACT

Finite element methods are applied to the problem of characterizing linear, anisotropic elastic solids. The conventional finite element displacement formulation is used to simulate an elastic material in plane stress. An inverted finite element formulation is then applied, and the characterizing six material constants are calculated as numerical results.

A possible test device for the experimental characterization of anisotropic solids is postulated, the precision of displacement measurements to be required for such a device being determined by random perturbation analysis. Numerical constants accurate to within three percent are predicted if a precision of one part in eight hundred ($1/800$) can be measured. Numerical constants accurate to within one percent are predicted if a precision of one part in eight thousand ($1/8000$) can be measured in the test device.

TABLE OF CONTENTS

I.	INTRODUCTION -----	6
II.	THEORETICAL BACKGROUND -----	8
	A. THEORY OF ELASTICITY -----	8
	1. Generalized Linear, Anisotropic Materials in Three-dimensions -----	8
	2. Elasticity in Plane Stress -----	10
	3. Elasticity in Plane Strain -----	12
	4. Characterization of Anisotropic Materials in Planar Elasticity -----	14
	B. FINITE ELEMENT FORMULATION -----	15
	1. Finite Element Geometry -----	15
	2. Displacement Formulation for the Finite Element -----	16
	3. Material Characterization Using Finite Elements -----	18
	C. CHARACTERIZATION USING FINITE ELEMENTS: AN EXAMPLE -----	18
III.	EXPERIMENTAL SIMULATION AND PROCEDURE -----	23
	A. NODAL DISPLACEMENTS -----	23
	1. Convergence Study -----	23
	B. COMPUTER PROGRAM DIASTIC -----	25
	1. Program Organization -----	28
	2. Eigenvalue Criterion of Matrix Singularity -----	30
	C. DISPLACEMENT PERTURBATIONS -----	30
IV.	RESULTS AND CONCLUSIONS -----	36
	A. PERTURBATION RESULTS -----	36
	B. CONCLUSIONS -----	38
APPENDIX A	Derivation of Stiffness Matrices for a 4-noded Rectangular Finite Element -----	40

APPENDIX B	Computer Program DPLISOP	-----	52
APPENDIX C	Computer Program DLASTIC	-----	73
BIBLIOGRAPHY		-----	90
INITIAL DISTRIBUTION LIST		-----	92
FORM DD 1473		-----	93

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I. INTRODUCTION

Jaensson and Sundström [11] used stress and strain approximations, developed by a finite element analysis of the microstructure of a presumed isotropic WC-Co alloy, as source data from which Young's modulus and Poisson's ratio were calculated as the characterizing constants for the alloy. A more direct application of the finite element method to the characterization of elastic solids is possible and can be applied to a wider range of materials. When finite element characterization is applied to isotropic solids an immediate advantage is secured in that Young's modulus and Poisson's ratio are obtained simultaneously. When the finite element characterization technique is applied to a general linear, anisotropic elastic material an overwhelming advantage is obtained in that no other method provides an experimental determination of all six anisotropic material constants.

The ability to characterize general linear, anisotropic elastic materials makes the design of a test device incorporating the finite element characterizing technique highly desirable. The application of this technique to linear elastic problems has been largely ignored, although an increasing use of the technique has been made to problems in the mechanical characterization of physically and kinematically non-linear materials [6, 7, 9, 12, 13, 15, 16].

The work reported in this study concerns a general class of linear, anisotropic materials. Its goal is to identify and determine preliminary design parameters applicable to the design of a testing device to be used in the characterization of linear, anisotropic solids. The finite

element technique employed in this characterization scheme has been observed [12] to demonstrate a sensitivity to the measurements of displacement data, and it is a specific goal of this study to identify such sensitivities and the degree of measurement precision required.

The study has been divided into three parts. In the first part the analytics of the finite element characterization are examined. Rather tedious derivations result in closed-form formulations, around which the characterization technique is built. The second part of the study uses a structural analysis computer code to simulate necessary test displacement data. The third and final part of the study required the writing of a computer program incorporating the finite element analytics. Perturbations of the test displacement data and consequent computer solutions to the characterization problem permitted a determination of the measurement precisions required for the possible test device.

II. THEORETICAL BACKGROUND

A. THEORY OF ELASTICITY

All structural materials exhibit in various degrees the property of elasticity: external forces loading a structure produce deformations of the structure, and if these forces do not exceed a limiting value the deformations disappear when the forces are removed. The theory of elasticity provides mathematical relations between forces and displacements acting in a structure. The action of forces lead to the definition of the stress tensor, and the geometric deformations lead to the definition of the strain tensor. A relation between these two tensors is called a constitutive law [8, 10, 14, 19]. It is this law, in its simplest formulation, which is paramount to the characterization of linear elastic solids.

1. Generalized Linear, Anisotropic Materials in Three-dimensions

The general form of the constitutive law in the theory of elasticity, a generalized Hooke's law, is the set of functions relating stresses to strains:

$$\begin{aligned}\sigma_x &= f_1(\epsilon_{xx}, \epsilon_{yy}, \epsilon_{zz}, \gamma_{xy}, \gamma_{xz}, \gamma_{yz}) \\ \sigma_y &= f_2(\epsilon_{xx}, \epsilon_{yy}, \epsilon_{zz}, \gamma_{xy}, \gamma_{xz}, \gamma_{yz}) \\ \sigma_z &= f_3(\epsilon_{xx}, \epsilon_{yy}, \epsilon_{zz}, \gamma_{xy}, \gamma_{xz}, \gamma_{yz}) \\ \tau_{xy} &= f_4(\epsilon_{xx}, \epsilon_{yy}, \epsilon_{zz}, \gamma_{xy}, \gamma_{xz}, \gamma_{yz}) \\ \tau_{xz} &= f_5(\epsilon_{xx}, \epsilon_{yy}, \epsilon_{zz}, \gamma_{xy}, \gamma_{xz}, \gamma_{yz}) \\ \tau_{yz} &= f_6(\epsilon_{xx}, \epsilon_{yy}, \epsilon_{zz}, \gamma_{xy}, \gamma_{xz}, \gamma_{yz})\end{aligned}\tag{1}$$

where x, y, z are the axes of the orthogonal, Cartesian coordinate system; $\sigma_x, \sigma_y, \sigma_z$ are the normal stresses; $\tau_{xy}, \tau_{xz}, \tau_{yz}$ are the

shear stresses; ϵ_{xx} , ϵ_{yy} , ϵ_{zz} are the normal strains; and γ_{xy} , γ_{xz} , γ_{yz} are the shear strains.

In the case of small deformations the simplest form of functions (1) is a linear set of equations. Denoting the coefficients applicable to this linear formulation by a_{mn} :

$$\begin{aligned}
 \sigma_x &= a_{11}\epsilon_{xx} + a_{12}\epsilon_{yy} + a_{13}\epsilon_{zz} + a_{14}\gamma_{xy} + a_{15}\gamma_{xz} + a_{16}\gamma_{yz} \\
 \sigma_y &= a_{21}\epsilon_{xx} + a_{22}\epsilon_{yy} + a_{23}\epsilon_{zz} + a_{24}\gamma_{xy} + a_{25}\gamma_{xz} + a_{26}\gamma_{yz} \\
 \sigma_z &= a_{31}\epsilon_{xx} + a_{32}\epsilon_{yy} + a_{33}\epsilon_{zz} + a_{34}\gamma_{xy} + a_{35}\gamma_{xz} + a_{36}\gamma_{yz} \\
 \tau_{xy} &= a_{41}\epsilon_{xx} + a_{42}\epsilon_{yy} + a_{43}\epsilon_{zz} + a_{44}\gamma_{xy} + a_{45}\gamma_{xz} + a_{46}\gamma_{yz} \\
 \tau_{xz} &= a_{51}\epsilon_{xx} + a_{52}\epsilon_{yy} + a_{53}\epsilon_{zz} + a_{54}\gamma_{xy} + a_{55}\gamma_{xz} + a_{56}\gamma_{yz} \\
 \tau_{yz} &= a_{61}\epsilon_{xx} + a_{62}\epsilon_{yy} + a_{63}\epsilon_{zz} + a_{64}\gamma_{xy} + a_{65}\gamma_{xz} + a_{66}\gamma_{yz}
 \end{aligned} \tag{1a}$$

Using matrix notation, the set of functions (1a) is written:

$$\{\sigma\} = [A]\{\epsilon\} \tag{1b}$$

or,

$$\{\epsilon\} = [B]\{\sigma\} \tag{1c}$$

where $\{\sigma\}$ and $\{\epsilon\}$ are the vectors whose elements are, respectively, the stress and strain components of the stress and strain tensors at a point,

$$\{\sigma\} = [\sigma_x, \sigma_y, \sigma_z, \tau_{xy}, \tau_{xz}, \tau_{yz}]^T \text{ and } \{\epsilon\} = [\epsilon_{xx}, \epsilon_{yy}, \epsilon_{zz}, \gamma_{xy}, \gamma_{xz}, \gamma_{yz}]^T;$$

and $[A]$ and $[B]$ are symmetric, 6×6 matrices of coefficients,

$$[A]^{-1} = [B], \text{ consisting of twenty-one different constants, } a_{mn}.$$

Restricting attention to elasticity in two-dimensions, two distinct problem types are identified: plane stress problems, requiring

$$\sigma_z = \tau_{xz} = \tau_{yz} = 0,$$

and plane strain problems, requiring

$$\epsilon_{zz} = \gamma_{xz} = \gamma_{yz} = 0.$$

2. Elasticity in Plane Stress

In the plane stress problem the in-plane stresses are defined as:

$$\{\sigma^P\} = [\sigma_x, \sigma_y, \tau_{xy}]^T.$$

Corresponding to these in-plane stresses are in-plane strains:

$$\{\epsilon^P\} = [\epsilon_{xx}, \epsilon_{yy}, \gamma_{xy}]^T.$$

Other strain components also exist. The out-of-plane strain components are defined as:

$$\{\epsilon^{OP}\} = [\epsilon_{zz}, \gamma_{xz}, \gamma_{yz}]^T,$$

where $\{\epsilon^{OP}\}$ is the vector of these out-of-plane strains.

The constitutive equation for the plane stress problem is, in terms of equation (1b):

$$\begin{Bmatrix} \sigma^P \\ 0 \end{Bmatrix} = \begin{bmatrix} A_{11}^* & A_{12}^* \\ A_{12}^{*T} & A_{22}^* \end{bmatrix} \begin{Bmatrix} \epsilon^P \\ \epsilon^{OP} \end{Bmatrix}, \quad (2)$$

where

$$A_{11}^* = \begin{bmatrix} a_{11} & a_{12} & a_{14} \\ a_{21} & a_{22} & a_{24} \\ a_{41} & a_{42} & a_{44} \end{bmatrix}; \quad A_{12}^* = \begin{bmatrix} a_{13} & a_{15} & a_{16} \\ a_{23} & a_{25} & a_{26} \\ a_{43} & a_{45} & a_{46} \end{bmatrix}; \quad A_{22}^* = \begin{bmatrix} a_{33} & a_{35} & a_{36} \\ a_{53} & a_{55} & a_{56} \\ a_{63} & a_{65} & a_{66} \end{bmatrix};$$

and the a_{mn} are the linear coefficients in equation (1b).

From equation (2):

$$\{\epsilon^{OP}\} = -[(A_{22}^*)^{-1}(A_{12}^*)^T]\{\epsilon^P\} \quad (2a)$$

and,

$$\begin{aligned} \{\sigma^P\} &= [A_{11}^*]\{\epsilon^P\} + [A_{12}^*]\{\epsilon^{OP}\} \\ &= [A_{11}^* - (A_{12}^*)(A_{22}^*)^{-1}(A_{12}^*)^T]\{\epsilon^P\}. \end{aligned} \quad (2b)$$

In the plane stress problem the stress components only are in-plane:

the state of strain is three-dimensional with the out-of-plane components given by equation (2a).

Alternately, equation (2b) is written:

$$\{\sigma^P\} = [A']\{\epsilon^P\} \quad (2c)$$

where $[A']$ is a symmetric, 3 x 3 matrix of constants:

$$[A'] = [A_{11}^* - (A_{12}^*)(A_{22}^*)^{-1}(A_{12}^*)^T],$$

the determination of which characterizes a linear, anisotropic material in plane stress.

For perfectly isotropic materials submitted to a state of plane stress:

$$\begin{aligned} A'_{11} &= \frac{E}{1-\nu} = A'_{22} \\ A'_{12} &= \nu A'_{11} \\ A'_{13} &= A'_{23} = 0 \\ A'_{33} &= \frac{1}{2}(1-\nu)A'_{11} \end{aligned} \quad (3)$$

where E and ν are Young's modulus and Poisson's ratio, respectively, and A'_{11} , A'_{12} , A'_{13} , A'_{22} , A'_{23} , and A'_{33} are the six elastic constants forming $[A']$.

If an isotropic material with $E = 29.50 \times 10^3$ ksi and $\nu = 0.287$ is postulated and if this material were subjected to an external load configuration equivalent to plane stress, the constitutive equation becomes:

$$\{\sigma^P\} = \begin{bmatrix} 32147.998 & 9226.476 & 0.0 \\ & 32147.998 & 0.0 \\ \text{(symmetric)} & & 11460.0 \end{bmatrix} \{\epsilon^P\} \quad (3a)$$

3. Elasticity in Plane Strain

In the plane strain problem, the in-plane strains are defined as:

$$\{\epsilon^P\} = [\epsilon_{xx}, \epsilon_{yy}, \gamma_{xy}]^T.$$

Corresponding to these in-plane strains are in-plane stresses:

$$\{\sigma^P\} = [\sigma_x, \sigma_y, \tau_{xy}]^T.$$

Other stress components also exist, and these out-of-plane stresses are defined as:

$$\{\sigma^{OP}\} = [\sigma_z, \tau_{xz}, \tau_{yz}]^T,$$

where $\{\sigma^{OP}\}$ is the vector of out-of-plane stresses.

The constitutive equation for the plane strain problem is, in terms of equation (1c):

$$\begin{Bmatrix} \epsilon^P \\ 0 \end{Bmatrix} = \begin{bmatrix} B_{11}^* & B_{12}^* \\ B_{12}^{*T} & B_{22}^* \end{bmatrix} \begin{Bmatrix} \sigma^P \\ \sigma^{OP} \end{Bmatrix} \quad (4)$$

where

$$B_{11}^* = \begin{bmatrix} b_{11} & b_{12} & b_{14} \\ b_{21} & b_{22} & b_{24} \\ b_{41} & b_{42} & b_{44} \end{bmatrix}; \quad B_{12}^* = \begin{bmatrix} b_{13} & b_{15} & b_{16} \\ b_{23} & b_{25} & b_{26} \\ b_{43} & b_{45} & b_{46} \end{bmatrix}; \quad B_{22}^* = \begin{bmatrix} b_{33} & b_{35} & b_{36} \\ b_{53} & b_{55} & b_{56} \\ b_{63} & b_{65} & b_{66} \end{bmatrix};$$

and the b_{mn} are the linear coefficients in equation (1c).

From equation (4):

$$\{\sigma^{OP}\} = - [(B_{22}^*)^{-1} (B_{12}^{*T})] \{\sigma^P\} \quad (4a)$$

and,

$$\begin{aligned} \{\epsilon^P\} &= [B_{11}^*] \{\sigma^P\} + [B_{12}^*] \{\sigma^{OP}\} \\ &= [B_{11}^* - (B_{12}^*) (B_{22}^*)^{-1} (B_{12}^{*T})] \{\sigma^P\}. \end{aligned} \quad (4b)$$

In the plane strain problem the strain components only are in-plane: the state of stress is three-dimensional with the out-of-plane components given by equation (4a).

Alternately, equation (4b) is written:

$$\{\epsilon^P\} = [B'] \{\sigma^P\} \quad (4c)$$

where $[B']$ is a symmetric, 3×3 matrix of constants:

$$[B'] = [B_{11}^* - (B_{12}^*) (B_{22}^*)^{-1} (B_{12}^{*T})],$$

the determination of which characterizes a linear, anisotropic material in plane strain.

For perfectly isotropic materials submitted to a state of plane strain:

$$B'_{11} = \frac{1-\nu^2}{E} = B'_{22}$$

$$B'_{12} = -\frac{\nu}{1-\nu} = B'_{11}$$

$$B'_{13} = B'_{23} = 0$$

$$B'_{33} = \frac{2}{1-\nu} B'_{11}$$

where E and ν are Young's modulus and Poisson's ratio, respectively, and B'_{11} , B'_{12} , B'_{13} , B'_{22} , B'_{23} , and B'_{33} are the six elastic constants forming $[B']$.

4. Characterization of Anisotropic Materials in Planar Elasticity

For all planar problems in elasticity it is generally observed that the two-dimensional behavior of the material is defined by a constitutive law of in-plane components: a vector of in-plane components is related to a second vector of in-plane components by a symmetric, 3×3 matrix of elastic constants. The two cases of planar elasticity defined by equation (2c) and equation (4c) are not mathematically equivalent: the matrix of material constants $[A']$ is not the inverse of the matrix of material constants $[B']$. Equations (2c) and (4c) do permit, however, the writing of a general constitutive law applicable to all planar elasticity problems:

$$\{\sigma^P\} = [D]\{\epsilon^P\} \quad (5)$$

where $[D]$ is a 3×3 , symmetric matrix equal to matrix $[A']$ for plane stress

problems and equal to $[B']^{-1}$ for plane strain problems. It is the form of equation (5) with which this study is concerned: a technique of measuring the six, independent elastic constants,

$$[D] = \begin{bmatrix} c_{11} & c_{12} & c_{13} \\ & c_{22} & c_{23} \\ \text{(symmetric)} & & c_{33} \end{bmatrix},$$

is the goal of this study in the characterization of linear, anisotropic material constants.

B. FINITE ELEMENT FORMULATION

Analysis of the basic unit of a structure is the first step in the total analysis of the structure. The basic part under consideration here is the finite element. Attention is focused on problems of planar stress, although the technique is expressed in a form generally applicable to any linear, anisotropic material. The finite elements employed are necessarily restricted to two-dimensional configurations.

1. Finite Element Geometry

A myriad of two-dimensional shapes have been used in finite element analyses [5, 21], the simplest being that of a triangle. Another very simple shape is that of a rectangle. Other shapes, from trapezoids through to curved shapes, are also available. For the purpose of the characterization of a linear, anisotropic material the simple rectangle has been selected primarily because of its inherent symmetries and ease of mathematical description.

2. Displacement Formulation for the Finite Element

The displacement formulation for a finite element is well documented [4, 5, 17, 18, 21]. Appendix A presents the details for the formulation of a 4-noded, linear element.

The elemental stiffness is given by:

$$[k]^e = \iint_s [B]^T [D] [B] ds , \quad (6)$$

where $[k]^e$ is defined as the elemental stiffness matrix, and the indicated integration is performed across the area of the element.

The equation of static equilibrium for an element can be written as:

$$\{Q\}^e = [k]^e \{\delta\}^e , \quad (7)$$

where $\{Q\}^e$ is the vector of forces acting at the elemental nodes.

Appendix A provides the specific derivation of the elemental stiffness matrix $[k]^e$ for a 4-noded rectangular finite element in closed-form. The resulting symmetric, 8×8 matrix is shown as Figure 7 of that appendix.

In conventional finite element analysis, elemental stiffness contributions are summed at common nodal points throughout the structure by the method of direct stiffness. Forces are summed throughout the structure, and it is observed that only forces external to the structure remain because of the nullifying effects of internal nodal forces. Resultant forces are of two kinds: forces due to the reactions of imposed physical boundary conditions and forces loading the structure. The equation of equilibrium for the structure is:

$$\{R\} = [K]\{\delta\} , \quad (8)$$

where $\{R\}$ is the vector of resultant forces acting at each node, generally equal to zero at internal nodes of the structure; $[K]$ is the master stiffness matrix; and $\{\delta\}$ is the vector of displacements containing the displacements at each node of the structure.

Equation (6) defines elemental stiffness as a function of the particular geometry of the element and the six elastic constants. Performing the matrix multiplication indicated by equation (7) and factoring-out the vector of elastic constants, a modified elemental stiffness matrix is defined [12, 13]:

$$\{Q\}^e = [k^*]^e \{C\} , \quad (9)$$

where $[k^*]$ is the modified elemental stiffness matrix and $\{C\}$ is the vector of elastic constants, C_{11} , C_{12} , C_{13} , C_{22} , C_{23} , and C_{33} . This modified elemental stiffness matrix is not square: for the 4-noded element the modified stiffness matrix is of size 8×6 .

The members of the modified elemental stiffness matrix are functions of the particular geometry of the element and the nodal displacements, u_i and v_i , of the element. Figure 8, Appendix A, presents the closed-form solution for the modified stiffness matrix for the 4-noded rectangular element. Such a closed-form solution for the modified stiffness matrix is notationally straight forward, but a good deal of tedious manipulation is hidden behind the symbols: closed-form solutions have not been successfully obtained for the elemental stiffness matrix or the modified elemental stiffness matrix for an 8-noded rectangular element because of repeated algebraic difficulties.

A simplified direct summation procedure is used to form the master modified stiffness matrix, and the forces are again summed across the structure:

$$\begin{aligned} \{R\} &= [K^*] \{C\} \\ n \times 1 \quad n \times 6 \quad 6 \times 1 \quad . \end{aligned} \quad (10)$$

3. Material Characterization Using Finite Elements

Figure 8, Appendix A, presents the 8×6 modified stiffness matrix derived for the 4-noded rectangular element. A structure composed of a number of these elements will exhibit a master modified stiffness matrix of order $n \times 6$. For the planar problems of elasticity two degrees of freedom exist at each node in the structure, and n equals twice the number of joints.

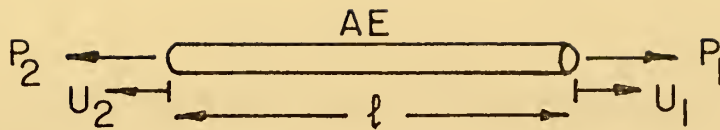
The over-determined system of equations (10) can be solved as follows:

$$\begin{aligned} [K^*]^T \{R\} &= [K^*]^T [K^*] \{C\} \\ ([K^*]^T [K^*])^{-1} [K^*]^T \{R\} &= \{C\} \quad , \end{aligned} \quad (11)$$

if the product $[K^*]^T [K^*]$ is non-singular.

C. CHARACTERIZATION USING FINITE ELEMENTS: AN EXAMPLE

Figure 1 presents an illustrative example of the use of finite elements in the characterization of elastic solids and is a modification after Kavanagh [12]. The problem is restricted to one-dimension and to one elastic constant, E .



TYPICAL FINITE ELEMENT, ONE-DIMENSIONAL.

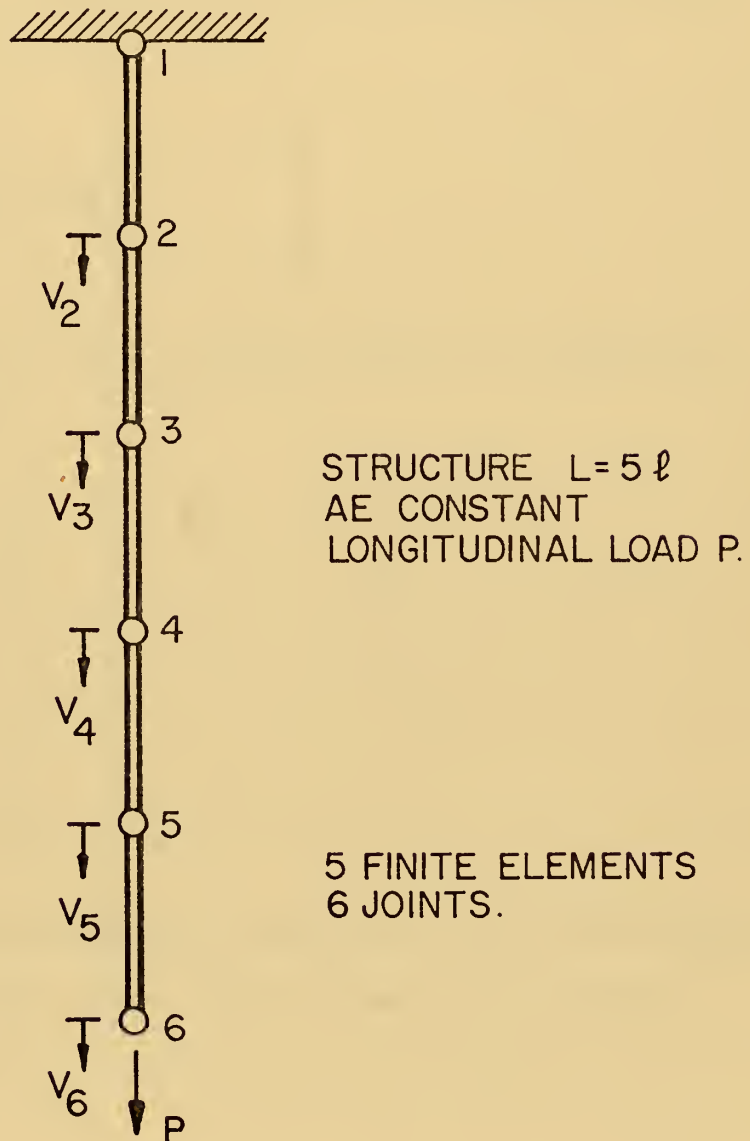


FIGURE 1 : ONE-DIMENSIONAL EXAMPLE PROBLEM
 ILLUSTRATING THE FINITE ELEMENT
 CHARACTERIZATION OF ELASTIC SOLIDS.

The example consists of a weight P , hanging from a slender rod of total length L and cross-sectional area A . The rod is divided into five finite elements, each element containing two nodes. There are a total of six structural joints, the first of which is constrained by a physically pinned connection.

The finite element in one-dimension yields an elemental stiffness matrix:

$$[k]^e = \frac{AE}{\ell} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} .$$

The master stiffness matrix for the overall structure is found by the direct stiffness method:

$$[K] = \frac{AE}{\ell} \begin{bmatrix} 1 & -1 & 0 & 0 & 0 & 0 \\ -1 & 2 & -1 & 0 & 0 & 0 \\ 0 & -1 & 2 & -1 & 0 & 0 \\ 0 & 0 & -1 & 2 & -1 & 0 \\ 0 & 0 & 0 & -1 & 2 & -1 \\ 0 & 0 & 0 & 0 & -1 & 1 \end{bmatrix} .$$

Imposition of the boundary condition at the first joint allows eliminating the first row and the first column, and the equation for structural equilibrium is written:

$$\begin{Bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ P \end{Bmatrix} = \frac{AE}{\ell} \begin{bmatrix} 2 & -1 & 0 & 0 & 0 \\ -1 & 2 & -1 & 0 & 0 \\ 0 & -1 & 2 & -1 & 0 \\ 0 & 0 & -1 & 2 & -1 \\ 0 & 0 & 0 & -1 & 1 \end{bmatrix} \begin{Bmatrix} v_2 \\ v_3 \\ v_4 \\ v_5 \\ v_6 \end{Bmatrix} .$$

Performing the indicated matrix multiplication and factoring the material constant, E:

$$\begin{Bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ P \end{Bmatrix} = \begin{bmatrix} (2v_2 - v_3) A/\ell \\ (-v_2 + 2v_3 - v_4) A/\ell \\ (-v_3 + 2v_4 - v_5) A/\ell \\ (-v_4 + 2v_5 - v_6) A/\ell \\ (-v_5 + v_6) A/\ell \end{bmatrix} \quad \begin{matrix} \{E\} \\ (1 \times 1) \end{matrix} .$$

$(5 \times 1) \qquad \qquad (5 \times 1)$

Setting $A = 1.0$, $\ell = 100.0$, and $P = 10.0$, assume that the measured deflections are $v_6 = 5.0$, $v_5 = 4.0$, $v_4 = 3.0$, $v_3 = 2.0$, $v_2 = 1.0$:

$$\begin{Bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 10.0 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0.01 \end{Bmatrix} \quad \{E\} .$$

Solving the equation for $\{E\}$:

$$[0 \ 0 \ 0 \ 0 \ 0.01] \begin{Bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 10.0 \end{Bmatrix} = [0 \ 0 \ 0 \ 0 \ 0.01] \begin{Bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0.01 \end{Bmatrix} \{E\}$$

$$0.10 = 0.0001 \{E\}$$

$$1000.0 = \{E\} \quad .$$

It is instructive to note that displacements calculated using this value of E would be identically those values assumed at joints 2, 3, 4, 5, and 6. In an actual experiment, however, the displacements obtained at the joints by measurement of the structure would be in error by some experimental amount. If E were known exactly, the E calculated would not be exact.

III. EXPERIMENTAL SIMULATION AND PROCEDURE

A. NODAL DISPLACEMENTS

The application of the finite element method to the characterization of elastic solids is predicated on the accurate knowledge of displacements at the nodes of the structure. Small errors in obtaining displacements will cause large errors to be introduced in the values of the calculated material constants, as reported by Kavanagh [12]. It has been necessary, therefore, to employ a simulation scheme to generate the displacements associated with a given loading configuration.

Appendix B consists of a listing of the digital computer program DPLISOP. DPLISOP is a double-precision version of PLISOP [3], a general purpose finite element program for planar problems in elasticity. DPLISOP was used to generate the nodal displacements required throughout this study.

1. Convergence Study

A convergence study using DPLISOP was performed. A steel plate, 10.0 inches square and 0.10 inches thick, was simulated in a state of plane stress. The top edge of the plate was constrained, and the bottom edge of the plate was uniformly loaded to a static load of 15.0 kips. Loading at the bottom edge nodes was simulated by the usual application of the consistent load vector concept [21]. Young's modulus and Poisson's ratio were taken to be $E = 29.50 \times 10^3$ ksi and $\nu = 0.287$, respectively. Finite element discretizations of one, four, and sixteen elements were used. Figure 2 shows these trial discretizations using the 4-noded element. Although DPLISOP is capable of utilizing 4-noded, 8-noded, and

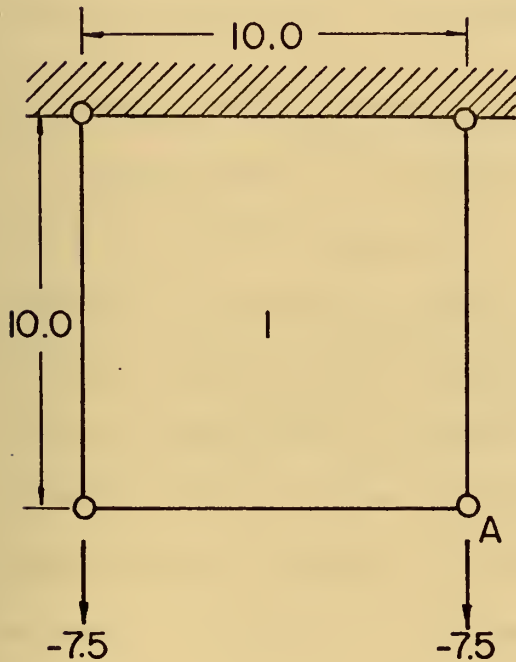
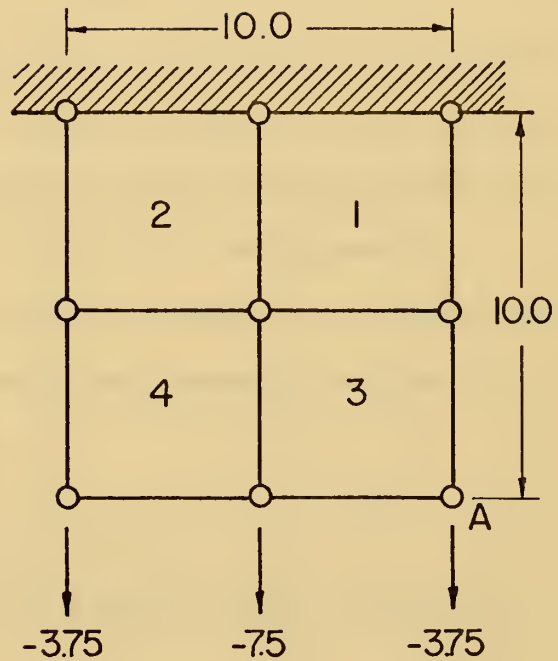
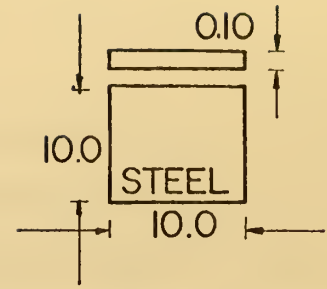
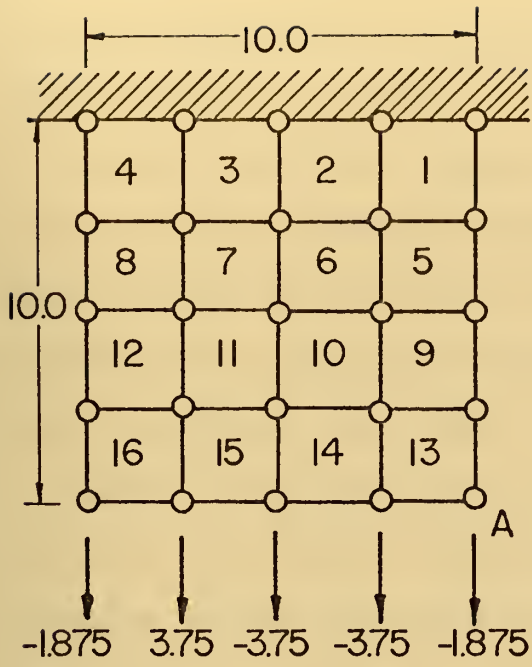


FIGURE 2: FINITE ELEMENT DISCRETIZATION FOR 4-NODED CONVERGENCE STUDY. STEEL PLATE, 10 BY 10, 0.10 THICK, TOTAL 15 KIP LOAD UNIFORMLY ALONG BOTTOM OF PLATE.

12-noded rectangular finite elements, runs were made for only the 4-noded and 8-noded elements.

Figure 3 presents the results of the convergence study for the 4-noded rectangular element, at Node A of the simulation structure. Similar results were obtained for the 8-noded rectangular element. The minimum number of elements necessary to ensure accurate displacements in both the horizontal and vertical directions was found to be four.

Implicit in the convergence study was the requirement that the experimental structure exhibit symmetries in both the horizontal and vertical directions and that elements be very nearly square to avoid geometrically induced anisotropic behavior. Figure 4 details the configuration, consisting of four, 4-noded rectangular finite elements, containing nine nodal points, and an overall 10.0 inches square geometry, selected for use in the remainder of the study.

B. COMPUTER PROGRAM DLASTIC

Program DLASTIC was written as a test program for the determination of the six elastic constants C_{11} , C_{12} , C_{13} , C_{22} , C_{23} , and C_{33} , by direct inversion of the finite element formulation (equations (10) and (11)), in order to characterize a general, linear anisotropic material. Double-precision displacements generated in DPLISOP were used as input data to DLASTIC and the six material constants calculated. These calculated numerical values proved to be exactly the characterizing six material constants of equation (3a). These results appear as a representative output which, with the listing of DLASTIC, form Appendix C.

Element geometry used in DLASTIC is restricted to the 4-noded rectangular finite element shape.

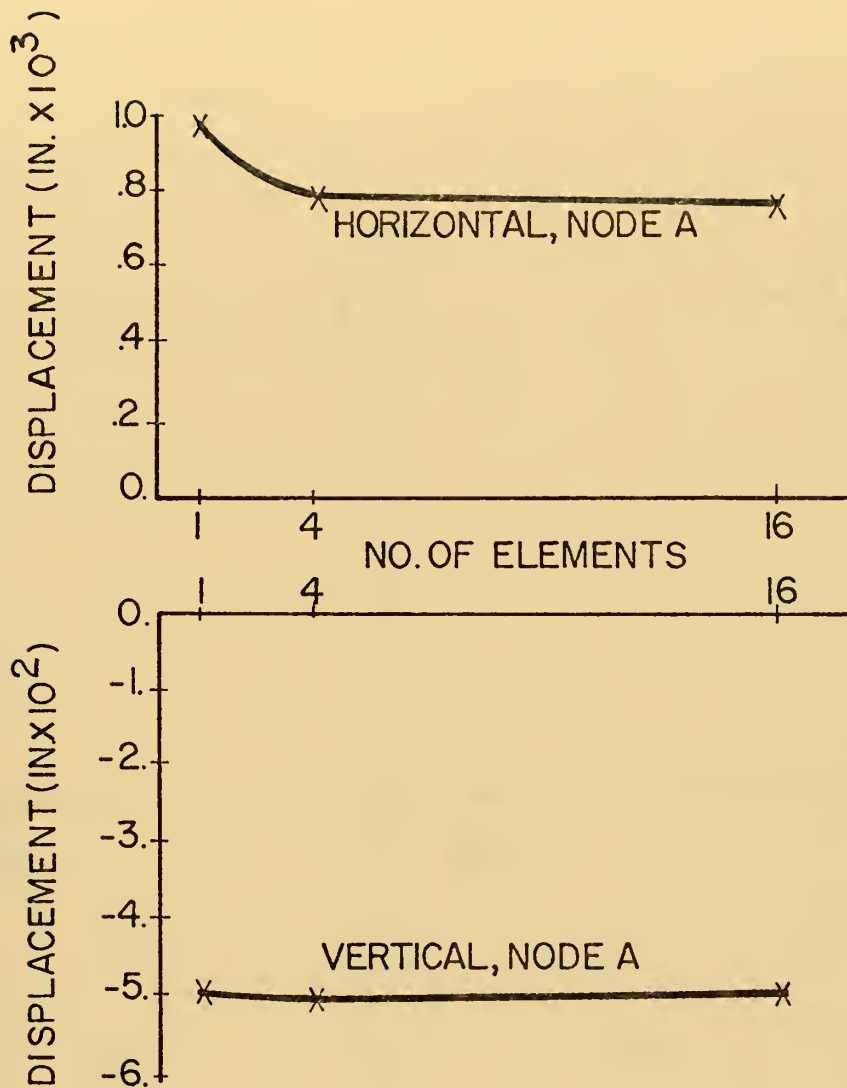


FIGURE 3: CONVERGENCE STUDY, NODE-A, USING 4-NODED RECTANGULAR FINITE ELEMENTS. STEEL PLATE IS 10 BY 10 SQUARE, 0.10 THICK, LOADED IN PLANE STRESS BY 15 KIP TOTAL LOAD UNIFORMLY DISTRIBUTED ALONG BOTTOM EDGE OF PLATE.

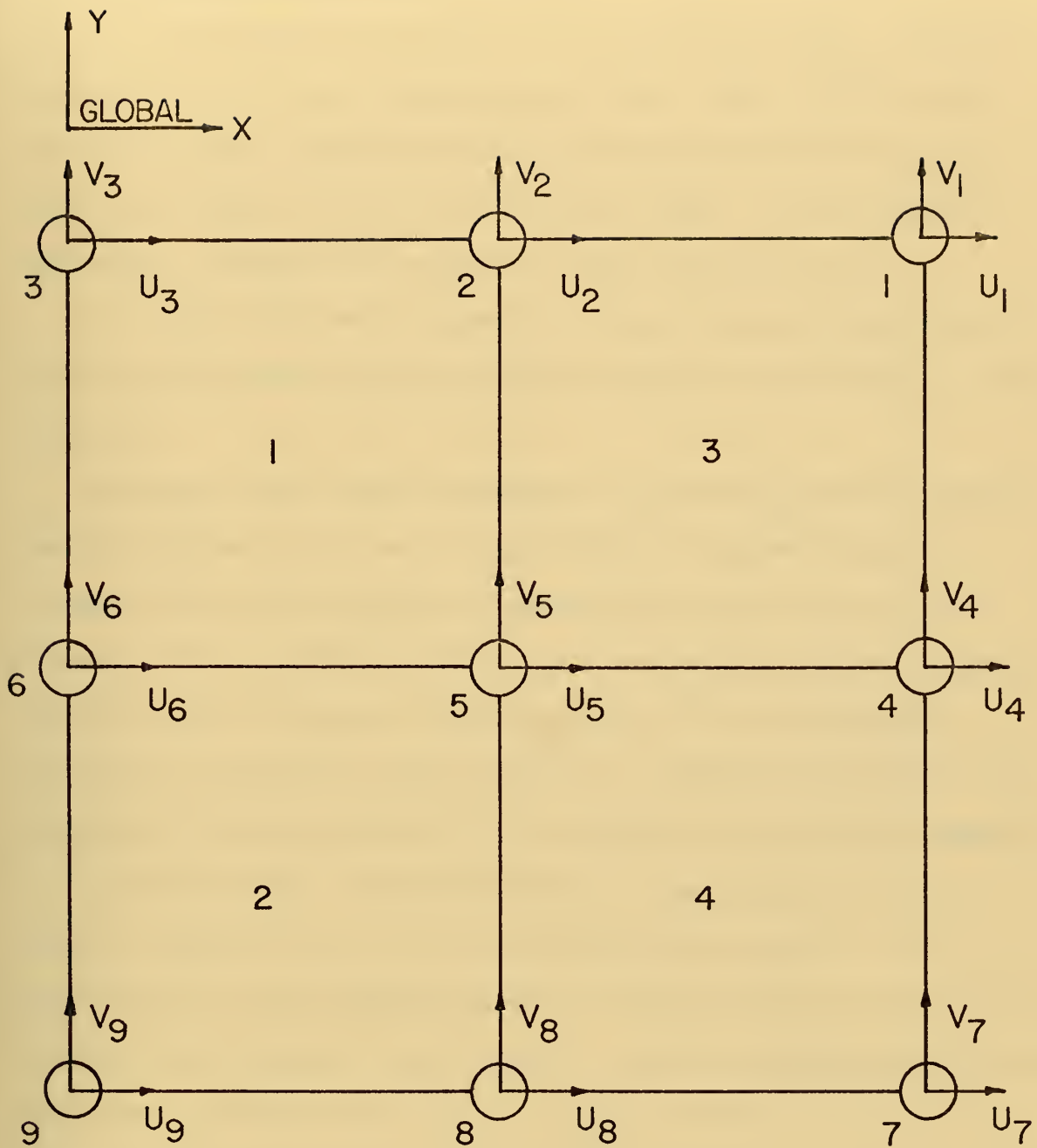


FIGURE 4 : 4-NODED 4-ELEMENT STRUCTURE SHOWING JOINT NUMBERING CONVENTION AND POSITIVE DISPLACEMENTS AT JOINTS.

1. Program Organization

Program DLASTIC was written as a modular algorithm: the MAIN program calls, in order, the subroutines INPUT, MERGE, BMULT, SINGL, INVRT, and ANSER. Subroutine MERGE, in turn, calls subroutine STIFF. Subroutine SINGL, in turn, calls subroutine EIGEN. Communication is maintained throughout the program by common blocks of storage.

The eight digit numbers appearing in the discussion which follows refer to the program sequencing numbers appearing in the program listing, Appendix C.

Subroutine INPUT (00000700-00001350) is called by the MAIN program. Punched data cards are read and echo-checked. Required input data includes structural parameter specifications, such as the number of elements and joints (NEL and NJT), nodal connectivity for every element (NCON), and the coordinates (COORD), forces (R), and displacements (U and V) for each joint in the structure. Specific directions for data preparation have been included within the MAIN program listing (Appendix C).

Subroutine MERGE (00001360-00001820) is next called by the MAIN program. MERGE forms the master modified stiffness matrix $[K^*]$, by superposition of the elemental modified stiffness matrices at appropriate common joints throughout the structure. Subroutine STIFF (00001830-00002910) is called in MERGE within a do-loop sequence indexed over the number of elements in the structure. Subroutine STIFF forms the elemental modified stiffness matrix.

Next called by the MAIN program is subroutine BMULT (00002920-00003270). In BMULT, the master modified stiffness matrix, BK, is premultiplied by

its transpose, BKT. The resulting symmetric matrix has been designated BSYM. The vector of nodal forces, R, is also premultiplied by BKT: this result has been designated BR.

The first multiplication in subroutine BMULT, the symmetric matrix BSYM, corresponds to the term $[K^*]^T [K^*]$ of equation (11). The MAIN program calls subroutine SINGL (00003280-00003530), which calls subroutine EIGEN (00003540-00005350) and a determination is made regarding the singularity or non-singularity of the matrix BSYM.

EIGEN calculates the six eigenvalues of BSYM associated with the as yet undetermined six material constants. SINGL causes the six eigenvalues to be printed as part of the program output, and inspection of this portion of the output is required by the program user to ascertain the validity of the total program output. Discussion of this requirement is deferred to the section following this discussion of program organization.

Subroutine INVRT (00005360-00006930) is next called from the MAIN program. The result of the first multiplication in subroutine BMULT, the matrix BSYM, is inverted using the Gauss-Jordan method. The determinant calculated in INVRT is ignored, for reasons which will become clear in the discussion of the eigenvalue criterion.

The vector of material constants, C, is calculated in subroutine ANSER (00006940-00007190) by premultiplying the vector BR, calculated in subroutine BMULT, by the inverted matrix BSYM, which occupies the storage location previously occupied by the non-inverted matrix. The resulting vector C, if valid, contains the six material constants C_{11} , C_{12} , C_{13} , C_{22} , C_{23} , and C_{33} . ANSER causes the constants to be printed as program output.

DLASTIC was written in FORTRAN IV and compiled with the G-compiler under release 18 on an IBM 360/67 machine. It required 84,000 bytes of core and used approximately one-second CPU time for a typical run.

2. Eigenvalue Criterion of Matrix Singularity

A matrix is said to be singular when the determinant of the matrix is equal to zero. However, when a matrix is ill-conditioned and nearly singular the numerical evaluation of the determinant is almost invariably completely unreliable. Round-off in most algorithms is sufficient to introduce an artificial round-off value of zero that invalidates the calculation of the determinant.

The calculation of the eigenvalues of a symmetric matrix is always a well conditioned numerical problem, even when the matrix is exactly singular. A better indication of singularity and near-singularity is obtained by examining all eigenvalues appropriate to the matrix. The condition number of the matrix is then obtained as the ratio of the largest absolute value of the eigenvalues to the smallest absolute value of the eigenvalues. If this condition number is greater than 10^{12} the matrix is too nearly singular to be inverted accurately.

In program DLASTIC it was necessary to employ the eigenvalue criterion of singularity. The symmetric matrix BSYM, formed in subroutine BMULT, was found to exhibit eigenvalues in the range 10^{-6} to 10^{-9} . The condition number is formed, 10^3 , and it is obvious that matrix BSYM is very far from any singularity.

C. DISPLACEMENT PERTURBATIONS

An exact characterization of an elastic solid was presented in a previous section (see the representative output, Appendix C). Such a

result utilized the full double-precision displacement data (sixteen digits) generated within the simulation program DPLISOP. Such exhaustive displacement data cannot be expected to be available from a physical experiment or possible test device. It was necessary, therefore, to perform a series of perturbation analyses in order to ascertain the lower limit of displacement data accuracy required for the characterization of elastic solids by the direct inversion of the finite element formulation.

A first set of truncated displacements was used to determine how sensitive the elastic constants would be to such approximations. Round-off to five, six, seven, and eight places of decimals was used (see Table 1). The elastic constants were calculated for each set of rounded displacements using program DLASTIC. The relative errors of the resulting material constants were then evaluated.

The results for the material constants and the associated errors are presented as Table 2.

Random perturbations, ranging in value from -5 to +5, were next applied to the rounded results in the fifth, sixth, seventh, and eighth places of decimals. Five different sets of displacements were obtained for each of the rounded results shown in Table 1.

Every set of displacements was used as input data for program DLASTIC, and all six material constants were obtained for each case. The results are shown in Tables 3 and 4.

JOINT	DISPLACEMENT	NON-RANDOM DISPLACEMENT PERTURBATIONS		
4 _U	-0.748329139132374D-03	-0.00075	-0.000748	-0.0007483
V	-0.252610022185014D-02	-0.00253	-0.002526	-0.0025261
5 _V	-0.235459058006389D-02	-0.00253	-0.002355	-0.0023546
6 _U	0.748329139132375D-03	0.00075	0.000748	0.0007483
V	-0.252610022185015D-02	-0.00253	-0.002526	-0.0025261
7 _U	-0.770070823585924D-03	-0.00077	-0.000770	-0.0007701
V	-0.500613331025850D-02	-0.00501	-0.005006	-0.0050061
8 _V	-0.497625861993874D-02	-0.00498	-0.004976	-0.0049763
9 _U	0.770070823585927D-03	0.00077	0.000770	0.0007701
V	-0.500613331025850D-02	-0.00501	-0.005006	-0.0050061

RUN1 RUN2 RUN3 RUN4

TABLE 1 :NON-RANDOM PERTURBATIONS (SIMPLE ROUNDING) OF NON-ZERO HORIZONTAL AND VERTICAL NODAL DISPLACEMENTS, SIMULATED 10.0BY10.0, 0.10 THICK STEEL STRUCTURE IN PLANE STRESS. 15.0 KIPS LOAD.

EXACT VALUES, ANISOTROPIC (PL. STRESS)					
C_{11}	C_{12}	C_{13}	C_{22}	C_{23}	C_{33}
32147.998	9226.476	0.0	32147.998	0.0	11460.0

RUN 1			RUN 2		
		ERROR, %			ERROR, %
C_{11}	32757.	1.9	C_{11}	32093.	0.2
C_{12}	9381.	1.7	C_{12}	9211.	0.2
C_{13}	0.	0.0	C_{13}	0.	0.0
C_{22}	32159.	0.0	C_{22}	32145.	0.0
C_{23}	0.	0.0	C_{23}	0.	0.0
C_{33}	11421.	0.3	C_{33}	11471.	0.0

RUN 3			RUN 4		
		ERROR, %			ERROR, %
C_{11}	32145.	0.0	C_{11}	32148.	0.0
C_{12}	9226.	0.0	C_{12}	9226.	0.0
C_{13}	0.	0.0	C_{13}	0.	0.0
C_{22}	32148.	0.0	C_{22}	32148.	0.0
C_{23}	0.	0.0	C_{23}	0.	0.0
C_{33}	11463.	0.0	C_{33}	11461.	0.0

TABLE 2: THE SIX ANISOTROPIC CONSTANTS CALCULATED IN PROGRAM PLASTIC, BY ROUNDING OF DISPLACEMENT DATA.

	EXACT	PERTURBATION 5				
		1	2	3	4	5
C_{11}	32147.998	39194.	35217.	31973.	29993.	36222.
C_{12}	9226.476	10978.	9954.	9243.	8993.	10353.
C_{13}	0.0	1041.	620.	328.	-759.	-414.
C_{22}	32147.998	32332.	32402.	32352.	32261.	32401.
C_{23}	0.0	95.	52.	-5.	-53.	-61.
C_{33}	11460.761	10665.	9882.	11924.	11314.	8591.

		PERTURBATION 6				
		Percent error				
C_{11}	32147.998	31290.	31660.	31553.	32201.	32136.
C_{12}	9226.476	9005.	9108.	9058.	9237.	9214.
C_{13}	0.0	32.	-105.	75.	-13.	-13.
C_{22}	32147.998	32108.	32115.	32130.	32174.	32136.
C_{23}	0.0	13.	-18.	6.	-7.	4.
C_{33}	11460.761	11594.	11518.	11699.	11481.	11618.

TABLE 3: THE SIX ANISOTROPIC MATERIAL CONSTANTS CALCULATED IN PROGRAM ELASTIC, BASED ON RANDOM PERTURBATIONS TO THE FIFTH AND SIXTH DECIMAL PLACES OF DISPLACEMENTS.

EXACT	PERTURBATION 7				
	1	2	3	4	5
C ₁₁	32153.	32102.	32074.	32094.	32103.
C ₁₂	9227.	9215.	9211.	9215.	9214.
C ₁₃	-6.	-3.	-9.	-14.	-4.
C ₂₂	32146.	32144.	32145.	32145.	32143.
C ₂₃	-2.	-2.	-1.	-2.	-1.
C ₃₃	11447.	11467.	11465.	11458.	11468.

EXACT	PERTURBATION 8				
	1	2	3	4	5
C ₁₁	32142.	32154.	32145.	32155.	32141.
C ₁₂	9225.	9228.	9226.	9228.	9225.
C ₁₃	0.	0.	0.	-1.	-1.
C ₂₂	32148.	32148.	32148.	32148.	32147.
C ₂₃	0.	0.	0.	0.	0.
C ₃₃	11463.	11459.	11459.	11462.	11463.

TABLE 4: THE SIX ANISOTROPIC MATERIAL CONSTANTS CALCULATED IN PROGRAM DLASTIC, BASED ON RANDOM PERTURBATIONS TO THE SEVENTH AND EIGHTH DECIMAL PLACES OF DISPLACEMENTS.

IV. RESULTS AND CONCLUSIONS

The six, anisotropic material constants appropriate to the plane stress elasticity problem and the particular material were calculated in equation (3a). Using the full double-precision displacement data generated by simulating an elastic structure in a state of plane stress program DLASTIC was used to calculate numerical values for these same six, anisotropic constants. Both sets of numerical constants were identical. For the purpose of discussion these values for the six material constants have been considered the exact numerical values of the constants.

Program DLASTIC required closed-form solutions to the stiffness matrices of the finite element formulation. Closed-form solutions for the elemental stiffness matrix, $[k]$, and the modified elemental stiffness matrix, $[k^*]$, were successfully obtained for the 4-noded, rectangular finite element. These results form Figures 7 and 8, Appendix A.

A. PERTURBATION RESULTS

Reasonable numerical values for the six elastic constants were obtained, based on approximate displacement data generated by rounding to as few as five decimal places. Table 2 presents the calculated values for these constants. Eigenvalues were also calculated: the four runs using simple roundings as approximations returned eigenvalues in the range 10^{-6} to 10^{-9} . Approximate displacement data to the fifth decimal place caused the returned numerical values of the constants to be within two percent of the exact numerical values, and approximate

displacement data to the sixth, seventh, and eighth decimal places resulted in numerical values within half-of-one percent of the exact values of the material constants.

Twenty sets of displacement data were generated for the random analyses. All twenty data sets were run in program DLASTIC and numerical values for the elastic constants calculated.

The numerical values for the calculated elastic constants are presented in Tables 3 and 4. The calculated eigenvalues for all random perturbations were within the range of 10^{-6} to 10^{-9} . Comparison of these calculated material constants to the exact values gives an idea of the sensitivity of the elastic constants to the accuracy of the displacement data.

Random perturbations applied to the results rounded to five decimal places gave the largest errors in the elastic constants. Worse values were consistently the calculated constant C_{13} , the values of which ranged from -759.5 ksi to 1041.0 ksi.

Simulated structural displacements accurate to 10^{-5} inches, the sixth decimal place of displacements being random, returned numerical values for the material constants generally less than three percent in error. Values of C_{13} , though not identically zero, were sufficiently small to consider them calculated zero except in two cases. Larger errors were experienced in the calculated values for C_{11} and C_{12} for all five sets.

Actual displacement measurements for a test device or physical experiment could be expected to require a horizontal displacement data precision of one part in eight hundred (1/800) and a vertical

displacement precision of one part in five hundred (1/500). These precisions could be expected to yield numerical values for the material constants in the range of two to three percent error.

Perturbations to the seventh and eighth decimal places returned material constants less than one percent in error (Table 4). An actual test device would require a horizontal measurement precision of one part in eight thousand (1/8000) and a vertical precision of one part in five thousand (1/5000), based on the simulated structural displacements being accurate to 10^{-6} inches and the seventh decimal place of displacements being random. The test device could be expected to yield numerical values for the material constants much less than one percent in error.

The numerical values returned from perturbations of displacement data to the eighth decimal place indicated required precisions of one part in eighty thousand (1/80000), horizontal, and one part in fifty thousand (1/50000), vertical. These precisions could be expected to yield numerical values for the material constants much less than one percent in error.

B. CONCLUSIONS

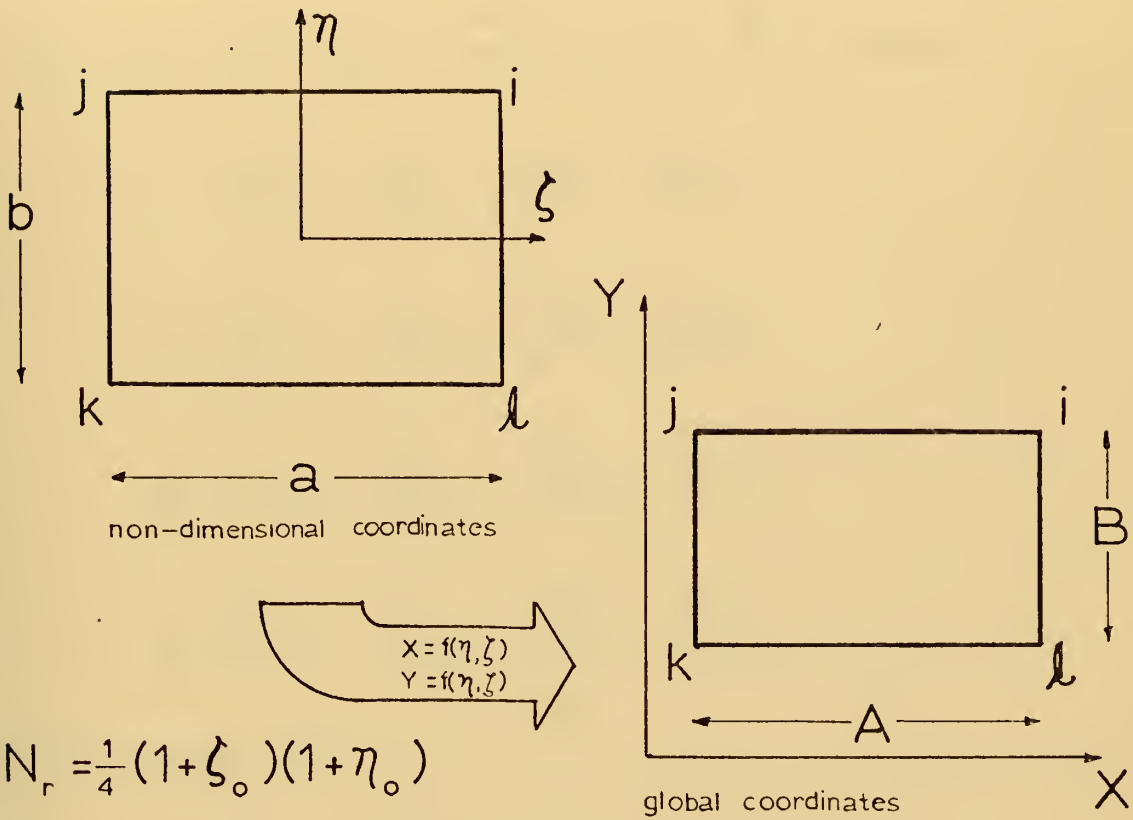
As a consequence of this study of the finite element technique, as utilized in the characterization of linear, anisotropic elastic solids, several conclusions can be drawn:

1. The finite element formulation under study can be applied to the mechanical characterization of linear, anisotropic solids.
2. The accuracy of characterization is dependent on the sensitivity of the finite element formulation.
3. The sensitivity of the finite element formulation is dependent on the precision with which displacement data is measured.

4. A device designed to determine the six linear, anisotropic elastic constants to an accuracy of three percent will require a precision in displacement measurements of at least one part in eight hundred ($1/800$).
5. A device designed to determine the six linear, anisotropic elastic constants to an accuracy of one percent will require a precision in displacement measurements of at least one part in eight thousand ($1/8000$).

APPENDIX A

Derivation of stiffness matrices for a 4-noded finite element



$$N_r = \frac{1}{4} (1 + \zeta_o)(1 + \eta_o)$$

$$\zeta_o = \zeta_r \zeta \quad \zeta_r = \pm 1$$

$$\eta_o = \eta_r \eta \quad \eta_r = \pm 1$$

$$N_i = \frac{1}{4} (1 + \zeta)(1 + \eta)$$

$$N_j = \frac{1}{4} (1 - \zeta)(1 + \eta)$$

$$N_k = \frac{1}{4} (1 - \zeta)(1 - \eta)$$

$$N_l = \frac{1}{4} (1 + \zeta)(1 - \eta)$$

$$x = N_i x_i + N_j x_j + N_k x_k + N_l x_l$$

$$y = N_i y_i + N_j y_j + N_k y_k + N_l y_l$$

$$\frac{\partial N_i}{\partial \zeta} = \frac{1}{4}(1+\eta) \quad \frac{\partial N_i}{\partial \eta} = \frac{1}{4}(1+\zeta) \quad \frac{\partial N_j}{\partial \zeta} = -\frac{1}{4}(1+\eta) \quad \frac{\partial N_j}{\partial \eta} = \frac{1}{4}(1-\zeta)$$

$$\frac{\partial N_k}{\partial \zeta} = \frac{1}{4}(\eta-1) \quad \frac{\partial N_k}{\partial \eta} = \frac{1}{4}(\zeta-1) \quad \frac{\partial N_\ell}{\partial \zeta} = \frac{1}{4}(1-\eta) \quad \frac{\partial N_\ell}{\partial \eta} = -\frac{1}{4}(1+\zeta)$$

$$\frac{\partial x}{\partial \zeta} = \frac{\partial N_i}{\partial \zeta} x_i + \frac{\partial N_j}{\partial \zeta} x_j + \frac{\partial N_k}{\partial \zeta} x_k + \frac{\partial N_\ell}{\partial \zeta} x_\ell$$

$$\frac{\partial x}{\partial \eta} = \frac{\partial N_i}{\partial \eta} x_i + \frac{\partial N_j}{\partial \eta} x_j + \frac{\partial N_k}{\partial \eta} x_k + \frac{\partial N_\ell}{\partial \eta} x_\ell$$

$$\frac{\partial y}{\partial \zeta} = \frac{\partial N_i}{\partial \zeta} y_i + \frac{\partial N_j}{\partial \zeta} y_j + \frac{\partial N_k}{\partial \zeta} y_k + \frac{\partial N_\ell}{\partial \zeta} y_\ell$$

$$\frac{\partial y}{\partial \eta} = \frac{\partial N_i}{\partial \eta} y_i + \frac{\partial N_j}{\partial \eta} y_j + \frac{\partial N_k}{\partial \eta} y_k + \frac{\partial N_\ell}{\partial \eta} y_\ell$$

Define the Jacobian, [J]:

$$[J(\eta, \zeta)] = \begin{bmatrix} \frac{\partial x}{\partial \zeta} & \frac{\partial y}{\partial \zeta} \\ \frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta} \end{bmatrix} = \begin{bmatrix} \frac{\partial N_i}{\partial \zeta} & \frac{\partial N_j}{\partial \zeta} & \frac{\partial N_k}{\partial \zeta} & \frac{\partial N_\ell}{\partial \zeta} \\ \frac{\partial N_i}{\partial \eta} & \frac{\partial N_j}{\partial \eta} & \frac{\partial N_k}{\partial \eta} & \frac{\partial N_\ell}{\partial \eta} \end{bmatrix} \begin{bmatrix} x_i & y_i \\ x_k & y_k \\ x_j & y_j \\ x_\ell & y_\ell \end{bmatrix}$$

$$= \frac{1}{2} \begin{bmatrix} A & 0 \\ 0 & B \end{bmatrix} \quad (a)$$

Define the displacement field, $\{\delta\} = \begin{Bmatrix} u \\ v \end{Bmatrix}$

$$u = N_i u_i + N_j u_j + N_k u_k + N_l u_l$$

$$v = N_i v_i + N_j v_j + N_k v_k + N_l v_l$$

(b)

$$\frac{\partial u}{\partial x} = \frac{\partial N_i}{\partial x} u_i + \frac{\partial N_j}{\partial x} u_j + \frac{\partial N_k}{\partial x} u_k + \frac{\partial N_l}{\partial x} u_l$$

$$\frac{\partial u}{\partial y} = \frac{\partial N_i}{\partial y} u_i + \frac{\partial N_j}{\partial y} u_j + \frac{\partial N_k}{\partial y} u_k + \frac{\partial N_l}{\partial y} u_l$$

$$\frac{\partial v}{\partial x} = \frac{\partial N_i}{\partial x} v_i + \frac{\partial N_j}{\partial x} v_j + \frac{\partial N_k}{\partial x} v_k + \frac{\partial N_l}{\partial x} v_l$$

$$\frac{\partial v}{\partial y} = \frac{\partial N_i}{\partial y} v_i + \frac{\partial N_j}{\partial y} v_j + \frac{\partial N_k}{\partial y} v_k + \frac{\partial N_l}{\partial y} v_l$$

In general,

$$\frac{\partial N_r}{\partial x} = \frac{\partial N_r}{\partial \zeta} \frac{\partial \zeta}{\partial x} + \frac{\partial N_r}{\partial \eta} \frac{\partial \eta}{\partial x}$$

$$\frac{\partial N_r}{\partial y} = \frac{\partial N_r}{\partial \zeta} \frac{\partial \zeta}{\partial y} + \frac{\partial N_r}{\partial \eta} \frac{\partial \eta}{\partial y}$$

or,

$$\frac{\partial N_r}{\partial \zeta} = \frac{\partial N_r}{\partial x} \frac{\partial x}{\partial \zeta} + \frac{\partial N_r}{\partial y} \frac{\partial y}{\partial \zeta}$$

(c)

$$\frac{\partial N_r}{\partial \eta} = \frac{\partial N_r}{\partial x} \frac{\partial x}{\partial \eta} + \frac{\partial N_r}{\partial y} \frac{\partial y}{\partial \eta}$$

Thus,

$$\begin{aligned}
 \begin{Bmatrix} \frac{\partial N_r}{\partial \zeta} \\ \frac{\partial N_r}{\partial \eta} \end{Bmatrix} &= \begin{bmatrix} \frac{\partial x}{\partial \zeta} & \frac{\partial y}{\partial \zeta} \\ \frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta} \end{bmatrix} \begin{Bmatrix} \frac{\partial N_r}{\partial x} \\ \frac{\partial N_r}{\partial y} \end{Bmatrix} \\
 &= [J(\eta, \zeta)] \begin{Bmatrix} \frac{\partial N_r}{\partial x} \\ \frac{\partial N_r}{\partial y} \end{Bmatrix} \\
 \begin{Bmatrix} \frac{\partial N_r}{\partial x} \\ \frac{\partial N_r}{\partial y} \end{Bmatrix} &= [J(\eta, \zeta)]^{-1} \begin{Bmatrix} \frac{\partial N_r}{\partial \zeta} \\ \frac{\partial N_r}{\partial \eta} \end{Bmatrix} \quad (d)
 \end{aligned}$$

where,

$$[J(\eta, \zeta)]^{-1} = 2 \begin{bmatrix} 1/A & 0 \\ 0 & 1/B \end{bmatrix} .$$

For each $r, r = (i, j, k, \ell)$:

$$\begin{aligned}
 \begin{Bmatrix} \frac{\partial N_i}{\partial x} \\ \frac{\partial N_i}{\partial y} \end{Bmatrix} &= \frac{1}{2} \begin{Bmatrix} \frac{1+\eta}{A} \\ \frac{1+\zeta}{B} \end{Bmatrix} \quad \begin{Bmatrix} \frac{\partial N_j}{\partial x} \\ \frac{\partial N_j}{\partial y} \end{Bmatrix} = \frac{1}{2} \begin{Bmatrix} \frac{-(1+\eta)}{A} \\ \frac{(1-\zeta)}{B} \end{Bmatrix} \quad (e)
 \end{aligned}$$

$$\begin{Bmatrix} \frac{\partial N_k}{\partial x} \\ \frac{\partial N_k}{\partial y} \end{Bmatrix} = \frac{1}{2} \begin{Bmatrix} \frac{(\eta-1)}{A} \\ \frac{(\zeta-1)}{B} \end{Bmatrix} \begin{Bmatrix} \frac{\partial N_\ell}{\partial x} \\ \frac{\partial N_\ell}{\partial y} \end{Bmatrix} = \frac{1}{2} \begin{Bmatrix} \frac{1-\eta}{A} \\ \frac{-(1+\zeta)}{B} \end{Bmatrix} \quad (e)$$

In general, for any two-dimensional element:

$$[B_r] = \begin{bmatrix} \frac{\partial N_r}{\partial x} & 0 \\ 0 & \frac{\partial N_r}{\partial y} \\ \frac{\partial N_r}{\partial y} & \frac{\partial N_r}{\partial x} \end{bmatrix} \quad (f)$$

and $[B] = [B_i \ B_j \ B_k \ B_\ell]$ for the 4-noded rectangular element.

Specifically:

$$[B] = \begin{bmatrix} \frac{\partial N_i}{\partial x} & 0 & \frac{\partial N_j}{\partial x} & 0 & \frac{\partial N_k}{\partial x} & 0 & \frac{\partial N_\ell}{\partial x} & 0 \\ 0 & \frac{\partial N_i}{\partial y} & 0 & \frac{\partial N_j}{\partial y} & 0 & \frac{\partial N_k}{\partial y} & 0 & \frac{\partial N_\ell}{\partial y} \\ \frac{\partial N_i}{\partial y} & \frac{\partial N_i}{\partial x} & \frac{\partial N_j}{\partial y} & \frac{\partial N_j}{\partial x} & \frac{\partial N_k}{\partial y} & \frac{\partial N_k}{\partial x} & \frac{\partial N_\ell}{\partial y} & \frac{\partial N_\ell}{\partial x} \end{bmatrix} \quad (g)$$

Substituting the functions from equation (e) into equation (g), the matrix of first partial derivatives of shape functions for a 4-noded rectangular finite element, [B], is formed, (Figure 5).

The matrix of elastic constants, for planar elasticity is:

$$[D] = \begin{bmatrix} C_{11} & C_{12} & C_{13} \\ & C_{22} & C_{23} \\ \text{(symmetric)} & & C_{33} \end{bmatrix} \quad (h)$$

where C_{11} , C_{12} , C_{13} , C_{22} , C_{23} , and C_{33} are the six elastic constants.

The products

$$[D][B], \quad \text{and}$$

$$[B]^T[D][B]$$

are then formed (Figure 6), where

$$\begin{aligned} d^2 &= \frac{(1+\eta)^2}{A^2} & h^2 &= \frac{(1-\zeta)^2}{B^2} & gf &= \frac{(1-\eta)(1+\zeta)}{AB} \\ g^2 &= \frac{(1+\zeta)^2}{B^2} & df &= \frac{(1-\eta)^2}{A^2} & gh &= \frac{(1-\zeta^2)}{B^2} \\ f^2 &= \frac{(1-\eta)^2}{A^2} & dg &= \frac{(1+\eta)(1+\zeta)}{AB} & fh &= \frac{(1-\eta)(1-\zeta)}{AB} \\ & & dh &= \frac{(1+\eta)(1-\zeta)}{AB} \end{aligned} \quad (i)$$

$$[B] = \frac{1}{2} \begin{bmatrix} \frac{1}{A}(1+\eta) & 0 & -\frac{1}{A}(1-\eta) & 0 & \frac{1}{A}(\eta-1) & 0 & \frac{1}{A}(1-\eta) & 0 \\ 0 & \frac{1}{B}(1+\xi) & 0 & \frac{1}{B}(1-\xi) & 0 & \frac{1}{B}(\xi-1) & 0 & -\frac{1}{B}(1+\xi) \\ \frac{1}{B}(1+\xi) & \frac{1}{A}(1+\eta) & \frac{1}{B}(1-\xi) & -\frac{1}{A}(1+\eta) & \frac{1}{B}(\xi-1) & \frac{1}{A}(\eta-1) & -\frac{1}{B}(1+\xi) & \frac{1}{A}(1-\eta) \end{bmatrix}$$

LET, $d = \frac{1}{A}(1+\eta)$

$f = \frac{1}{A}(1-\eta)$

$g = \frac{1}{B}(1+\xi)$

$h = \frac{1}{B}(1-\xi)$

THEN,

$$[B] = \frac{1}{2} \begin{bmatrix} d & 0 & -d & 0 & -f & 0 & f & 0 \\ 0 & g & 0 & h & 0 & -h & 0 & -g \\ g & d & h & -d & -h & -f & -g & f \end{bmatrix}$$

FIGURE 5: [B], THE MATRIX OF FIRST-PARTIAL DERIVATIVES OF SHAPE FUNCTIONS FOR A 4-NODED RECTANGULAR FINITE ELEMENT.

$$[D] [B] = \frac{1}{2} \begin{bmatrix} (c_{11}d + c_{13}g) & (c_{12}g + c_{13}d) & (-c_{11}f - c_{13}h) & (-c_{12}h - c_{13}f) & (c_{11}f - c_{13}g) & (-c_{12}g + c_{13}f) \\ (c_{12}d + c_{13}g) & (c_{22}g + c_{23}d) & (-c_{12}f - c_{23}h) & (-c_{22}h - c_{23}f) & (c_{12}f - c_{23}g) & (-c_{22}g + c_{23}f) \\ (c_{13}d + c_{23}g) & (c_{23}g + c_{33}d) & (-c_{13}f - c_{33}h) & (-c_{23}h - c_{33}f) & (c_{13}f - c_{33}g) & (-c_{23}g + c_{33}f) \end{bmatrix}$$

AND,

$$[B]^T [D] [B] = \frac{1}{4} \begin{bmatrix} (c_{11}d^2 + c_{13}gd + c_{13}gd + c_{33}g^2) & (c_{12}gd + c_{13}d^2) & (-c_{11}d^2 + c_{13}hd) & (c_{12}hd - c_{13}d^2) & (-c_{11}fd - c_{13}hd) & (-c_{12}hd - c_{13}fd) & (c_{11}fd - c_{13}gd) & (-c_{12}gd + c_{13}fd) & (-c_{13}gd + c_{33}g^2) & (-c_{23}g + c_{33}fg) \\ (c_{22}g^2 + c_{23}dg + c_{23}gd + c_{33}d^2) & (c_{22}hg - c_{23}dg) & (-c_{12}fg - c_{23}hg) & (c_{22}hg - c_{23}dg) & (-c_{12}fg - c_{23}hg) & (-c_{22}hg - c_{23}dg) & (c_{12}fg - c_{23}gd) & (-c_{22}g^2 + c_{23}gf) & (-c_{23}gd + c_{33}fd) & (-c_{23}gd + c_{33}fd) \\ (c_{11}d^2 + c_{13}hd - c_{13}dh + c_{33}h^2) & (-c_{12}d^2 + c_{13}hg) & (-c_{11}fd + c_{13}hd) & (-c_{12}hd - c_{13}d^2) & (-c_{11}fd - c_{13}hd) & (-c_{12}hd - c_{13}d^2) & (-c_{11}fd + c_{13}gd) & (c_{12}gd - c_{13}fd) & (-c_{13}fd + c_{33}gh) & (-c_{23}gh + c_{33}fh) \\ (c_{22}h^2 - c_{23}dh - c_{23}dh + c_{33}d^2) & (c_{22}h^2 - c_{23}dh) & (-c_{12}fh - c_{23}h^2) & (c_{22}h^2 - c_{23}dh) & (-c_{12}fh - c_{23}h^2) & (-c_{22}h^2 - c_{23}dh) & (c_{12}fh - c_{23}gd) & (-c_{22}gh + c_{23}fh) & (-c_{23}gd - c_{33}fd) & (-c_{23}gd - c_{33}fd) \\ (c_{11}f^2 + c_{13}hf + c_{13}fh + c_{33}h^2) & (-c_{12}fh - c_{23}h^2) & (-c_{11}f^2 + c_{13}hf) & (-c_{12}h^2 - c_{23}fh) & (-c_{11}f^2 + c_{13}hf) & (-c_{12}h^2 - c_{23}fh) & (-c_{11}f^2 + c_{13}gf) & (c_{12}gf - c_{13}f^2) & (-c_{13}fh + c_{33}gh) & (-c_{23}gh - c_{33}fh) \\ (c_{22}h + c_{23}fh + c_{23}hf + c_{33}f^2) & (c_{22}h^2 - c_{23}dh) & (c_{22}h + c_{23}fh) & (c_{22}h^2 - c_{23}dh) & (c_{22}h + c_{23}fh) & (c_{22}h^2 - c_{23}dh) & (-c_{12}fh + c_{23}gh) & (c_{22}gh - c_{23}fh) & (-c_{13}f^2 + c_{33}gf) & (-c_{23}gf - c_{33}f^2) \\ (c_{11}f^2 + c_{13}gf - c_{13}fg + c_{33}g^2) & (-c_{12}fg - c_{23}gd) & (-c_{13}fg + c_{33}gh) & (-c_{12}gd + c_{13}fd) & (-c_{13}fg + c_{33}gh) & (-c_{12}gd + c_{13}fd) & (c_{11}f^2 + c_{13}gf) & (c_{12}gf + c_{13}f^2) & (-c_{13}fg + c_{33}g^2) & (-c_{23}g^2 - c_{23}fg) & (-c_{23}gf + c_{33}f^2) \end{bmatrix}$$

(SYMMETRIC)

FIGURE 6

The elemental stiffness is defined:

$$[k]^e = \int_y \int_x [B]^T [D] [B] dx dy . \quad (j)$$

The product $[B]^T [D] [B]$, Figure 6, and equation (i) are functions in (η, ζ) . With $dx dy = \det [J(\eta, \zeta)]$:

$$[k]^e = \int_{-1}^1 \int_{-1}^1 [B]^T [D] [B] \det [J(\eta, \zeta)] d\eta d\zeta \quad (k)$$

where $\det [J(\eta, \zeta)] = \frac{1}{4} AB$ and the indicated integration is performed in the non-dimensionalized coordinate system. Performing this integration, the elemental stiffness matrix results (Figure 7).

The equation of static equilibrium for the element is written:

$$\{Q\}^e = [k]^e \{\delta\}^e \quad (l)$$

where $\{Q\}^e$ is the vector of forces acting at the nodes of the element; $[k]^e$ is the elemental stiffness matrix; and $\{\delta\}^e$ is the vector of elemental displacements,

$$\{\delta\} = \begin{Bmatrix} \delta_i \\ \delta_j \\ \delta_k \\ \delta_l \end{Bmatrix} = \begin{Bmatrix} u_i \\ v_i \\ u_j \\ v_j \\ u_k \\ v_k \\ u_l \\ v_l \end{Bmatrix} \quad (m)$$

(SYMMETRIC)

FIGURE 7: ELEMENTAL STIFFNESS MATRIX $[K]^e$ FOR A 4-NODED RECTANGULAR FINITE ELEMENT. CLOSED-FORM SOLUTION IN TERMS OF ELEMENT DIMENSIONS (A, B) AND THE VECTOR OF LINEAR, ANISOTROPIC MATERIAL CONSTANTS ($C_{11}, C_{12}, C_{13}, C_{22}, C_{23}, C_{33}$) MATRIX $[K]$ IS SYMMETRIC.

Performing the multiplication indicated by equation (l) and factoring the six material constants:

$$\{Q\}^e = [k^*]^e \{C\} \quad (n)$$

where $[k^*]^e$ is the modified elemental stiffness matrix, in terms of (m), (see Figure 8), and $\{C\}$ is the vector of material constants of equation (h).

$\frac{1}{3} \frac{B}{A} (-u_i - u_j - \frac{1}{2} u_k + \frac{1}{2} u_l)$	$\frac{1}{4} (v_i + v_j - v_k - v_l)$	$\frac{1}{2} (u_i - u_k)$ $+\frac{1}{3} \frac{B}{A} (v_i - v_j - \frac{1}{2} v_k + \frac{1}{2} v_l)$	0	$\frac{1}{3} \frac{A}{B} (v_i + \frac{1}{2} v_j - \frac{1}{2} v_k - v_l)$	$\frac{1}{4} (v_i - v_j - v_k + v_l)$ $+\frac{1}{3} \frac{A}{B} (u_i + \frac{1}{2} u_j - \frac{1}{2} u_k - u_l)$
0	$\frac{1}{4} (u_i - u_j - u_k + u_l)$	$\frac{1}{3} \frac{B}{A} (u_i - u_j - \frac{1}{2} u_k + \frac{1}{2} u_l)$	$\frac{1}{3} \frac{A}{B} (v_i + \frac{1}{2} v_j - \frac{1}{2} v_k - v_l)$	$\frac{1}{2} (v_i - v_k)$ $+\frac{1}{3} \frac{A}{B} (u_i + \frac{1}{2} u_j - \frac{1}{2} u_k - u_l)$	$\frac{1}{4} (u_i + u_j - u_k - u_l)$ $+\frac{1}{3} \frac{B}{A} (v_i - v_j - \frac{1}{2} v_k + \frac{1}{2} v_l)$
$\frac{1}{3} \frac{B}{A} (-u_i + u_j + \frac{1}{2} u_k - \frac{1}{2} u_l)$	$\frac{1}{4} (-v_i - v_j + v_k + v_l)$	$\frac{1}{2} (-u_j + u_l)$ $+\frac{1}{3} \frac{B}{A} (-v_i + v_j + \frac{1}{2} v_k - \frac{1}{2} v_l)$	0	$\frac{1}{3} \frac{A}{B} (\frac{1}{2} v_i + v_j - v_k - \frac{1}{2} v_l)$	$\frac{1}{4} (v_i - v_j - v_k + v_l)$ $+\frac{1}{3} \frac{A}{B} (-\frac{1}{2} u_i + u_j - \frac{1}{2} u_k - \frac{1}{2} u_l)$
0	$\frac{1}{4} (u_i - u_j - u_k + u_l)$	$\frac{1}{3} \frac{B}{A} (-u_i + u_j + \frac{1}{2} u_k - \frac{1}{2} u_l)$	$\frac{1}{3} \frac{A}{B} (\frac{1}{2} v_i - v_j - v_k - \frac{1}{2} v_l)$	$\frac{1}{2} (-v_i + v_l)$ $+\frac{1}{3} \frac{A}{B} (-\frac{1}{2} u_i + u_j - \frac{1}{2} u_k - \frac{1}{2} u_l)$	$\frac{1}{4} (-u_i - u_j + u_k + u_l)$ $+\frac{1}{3} \frac{B}{A} (-v_i + v_j + \frac{1}{2} v_k - \frac{1}{2} v_l)$
$\frac{1}{3} \frac{B}{A} (\frac{1}{2} u_i + \frac{1}{2} u_j + \frac{1}{2} u_k - u_l)$	$\frac{1}{4} (-v_i - v_j + v_k + v_l)$	$\frac{1}{2} (-u_i + u_k)$ $+\frac{1}{3} \frac{B}{A} (-\frac{1}{2} v_i + \frac{1}{2} v_j + v_k - v_l)$	0	$\frac{1}{3} \frac{A}{B} (-\frac{1}{2} v_i - v_j + v_k + \frac{1}{2} v_l)$	$\frac{1}{4} (-v_i + v_j + v_k - v_l)$ $+\frac{1}{3} \frac{A}{B} (-\frac{1}{2} u_i - u_j + u_k + \frac{1}{2} u_l)$
0	$\frac{1}{4} (-u_i + u_j + u_k - u_l)$	$\frac{1}{3} \frac{B}{A} (-\frac{1}{2} u_i + \frac{1}{2} u_j + u_k - u_l)$	$\frac{1}{3} \frac{A}{B} (-\frac{1}{2} v_i - v_j + v_k + \frac{1}{2} v_l)$	$\frac{1}{2} (-v_i + v_k)$ $+\frac{1}{3} \frac{A}{B} (-\frac{1}{2} u_i - u_j + u_k + \frac{1}{2} u_l)$	$\frac{1}{4} (-u_i - u_j + u_k + u_l)$ $+\frac{1}{3} \frac{B}{A} (-v_i + v_j + v_k - v_l)$
$\frac{1}{3} \frac{B}{A} (\frac{1}{2} u_i - \frac{1}{2} u_j - \frac{1}{2} u_k + u_l)$	$\frac{1}{4} (v_i + v_j - v_k - v_l)$	$\frac{1}{2} (u_j - u_l)$ $+\frac{1}{3} \frac{B}{A} (\frac{1}{2} v_i - \frac{1}{2} v_j - v_k + v_l)$	0	$\frac{1}{3} \frac{A}{B} (-v_i - \frac{1}{2} v_j + \frac{1}{2} v_k + v_l)$	$\frac{1}{4} (u_i + u_j - u_k - u_l)$ $+\frac{1}{3} \frac{A}{B} (\frac{1}{2} v_i - \frac{1}{2} v_j - \frac{1}{2} v_k + v_l)$
0	$\frac{1}{4} (-u_i + u_j + u_k - u_l)$	$\frac{1}{3} \frac{B}{A} (\frac{1}{2} u_i - \frac{1}{2} u_j - \frac{1}{2} u_k + u_l)$	$\frac{1}{3} \frac{A}{B} (-v_i - \frac{1}{2} v_j + \frac{1}{2} v_k + v_l)$	$\frac{1}{2} (v_j - v_l)$ $+\frac{1}{3} \frac{A}{B} (-u_i - \frac{1}{2} u_j + \frac{1}{2} u_k + u_l)$	$\frac{1}{4} (u_i + u_j - u_k - u_l)$ $+\frac{1}{3} \frac{B}{A} (\frac{1}{2} v_i - \frac{1}{2} v_j - \frac{1}{2} v_k + v_l)$

FIGURE 8 : $[k]^e$, MODIFIED ELEMENTAL STIFFNESS MATRIX FOR A 4-NODED RECTANGULAR FINITE ELEMENT, CLOSED-FORM SOLUTION IN TERMS OF ELEMENT DIMENSIONS (A,B) AND NODAL DISPLACEMENTS (u,v).

COMPUTER PROGRAM DPLISOP

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CHARACTERIZATION OF ELASTIC SOLIDS USING FINITE ELEMENT METHODS

U. S. NAVAL POSTGRADUATE SCHOOL
MONTEREY, CALIFORNIA
DECEMBER 1972

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*****
THIS PROGRAM PERFORMS A PLANE STRESS OR A PLANE STRAIN ANALYSIS
BY THE FINITE ELEMENT METHOD. NUMERICALLY INTEGRATED ISOPARAMETRIC
ELEMENTS ARE USED THROUGHOUT.
PROFESSOR GILLES CANTIN, MARCH 1972
** IMPLICIT REAL*8 (A-H,O-Z)
COMMON/SOL/BGK(312,64),ALOAD(312),DISP(312),ABGN
COMMON/INT/NEL,NJT,NMAT,NCLOAD,NPBC,NCON(100,14),NBC(156,2),NSTRES
1,NGP,LM(22),IJT(22)
COMMON/FLPL/BOARD(156,2),CLOAD(156,2),ELCON(10,2),TITLE(20),T

INPUT CARDS NEEDED

TITLE          (20A4)          ONE CARD PER PROBLEM
PROBLEM        (715,10X,1G24.16) ONE CARD PER PROBLEM
COL.           VARIABLE
1-5            NEL          NUMBER OF ELEMENTS (MAX. IS 100)
6-10           NJT          NUMBER OF JOINTS (MAX. IS 156)
11-15          NMAT         NUMBER OF MATERIALS (MAX. IS 10)
16-20          NCLOAD       NUMBER OF CONCENTRATED LOADS
21-25          NPBC         NUMBER OF JOINTS WITH BOUNDARY CONDITIONS
26-30          NSTRESS      ZERO FOR PLANE STRESS DESIRED IN THE INTEGRATION
31-35          NGP          NUMBER OF GAUSS POINTS T=0.0 OR BLANK IS SAME AS
46-70          T            SPECIMEN THICKNESS: (1G24.16)
UNIT           THICKNESS (T=1.0)

JOINT CARDS (CARD ONE:1110,2G31.16 CARD TWO:2G31.16) NJT PER PROB.
TWO CARDS PER JOINT ARE REQUIRED:
CARD ONE---JOINT IDENTIFICATION AND COORDINATES
1-10          JOINT NUMBER
11-41         X COORDINATE OF THE JOINT
42-72         Y COORDINATE OF THE JOINT
CARD TWO---LOADING AT THE JOINT
1-31         X COMPONENT OF THE LOAD AT THE JOINT
32-62         Y COMPONENT OF THE LOAD AT THE JOINT

```



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C MATERIAL PROPERTIES (1110,2625.16) NMAT CARDS PER PROBLEM
C COL. VARIABLE
C 1-10 MATERIAL NUMBER
C 11-35 YOUNG'S MODULUS
C 36-63 POISSON'S RATIO
C DIMENSION REACT(157,2),NCODE(156)
C EQUIVALENCE (BGK(1,1),REACT(1,1))
C
C ELEMENT CARDS (1514)
C COL. VARIABLE IDENTIFICATION NUMBER
C 1-4 ELEMENT CONNECTIVITY (COUNTERCLOCKWISE)
C 5-52 KIND 1 FOR FOUR NODAL POINT ELEMENT
C 53-56 KIND 2 FOR EIGHT NODAL POINT ELEMENT
C 57-60 TYPE 3 FOR TWELVE NODAL POINT ELEMENT
C THIS IS THE MATERIAL IDENTIFICATION NUMBER
C
C BOUNDARY CONDITION CARDS (215)
C COL. VARIABLE NPBC CARDS PER PROBLEM
C 1-5 JOINT NUMBER
C 6-10 CODE 1 FOR A ROLLER ON AN X AXIS, DISP. IN Y DIRECT. IS ZERO
C 2 FOR A ROLLER ON AN Y AXIS, DISP. IN X DIRECT. IS ZERO
C 3 FOR A FIXED JOINT, BOTH DISP. COMPONENTS ARE ZERO
C 4 FOR A ROLLER ON AN X AXIS, THE DISP. IN THE Y DIRECT.
C 5 SPECIFIED ON THE SECOND JOINT CARD IN COLS. 32-62
C 6 SPECIFIED ON THE SECOND JOINT CARD IN COLS. 32-62
C 7 SPECIFIED ON THE SECOND JOINT CARD IN COLS. 1-31
C DIRECTIONS AND PUNCHED IN COLS. 1-31 AND Y
C THE SECOND JOINT CARD (JOINT LOADING)
C 7 FOR A ROLLER ON AN X AXIS INCLINED ALPHA DEGREES
C WITH RESPECT TO THE X AXIS. THE JOINT CARD MAY
C CONTAIN A TANGENTIAL LOAD IN COL. 1-31 AND THE
C ANGLE ALPHA DEGREES IN COL. 32-62, SECOND JOINT CARD OF
C AS MANY PROBLEMS AS ONE DESIRES STOP IN THE FIRST FOUR COLUMNS
C A RUN MUST CONTAIN THE WORD
C PI=3.141592653589793D0
C 10 NBAND=0
C DO 20 I=1,156
C 20 NCODE(I)=0
C CALL INPUT
C NEQ=2*NJUT
C NPEL=NCON(1,13)*4
C DO 300 I=1,NPEL
C DO 100 J=1,NPEL
C 100 LM(J)=NCON(I,J)
C NPELM=NPEL-1
C DO 200 K=1,NPELM

```



```

200 JK=K+1
300 DO 200 L=JK,NPEL
    NBAND=MAX0(NBAND,IABS(LM(K)-LM(L)))
    CONTINUE
    NBAND=(NBAND+1)*2
    DO 350 I=1,NEQ
        ALOAD(I)=0.0D0
    DO 350 J=1,NBAND
        BGK(I,J)=0.0D0
        CALL MERGE(NPEL,NEQ)
        CALL BCOND(NBAND)
        CALL CHOL(NEQ,NBAND)
    DO 500 I=1,NP3C
        IJT=NBC(I,I)
        NCODE(IJT)=NBC(I,2)
    DO 600 I=1,NJT
        II=I*2-1
        KCDE=NCODE(I)
        IF(KCDE.EQ.7) GO TO 560
        IF((KCDE.EQ.5).OR.(KCDE.EQ.6)) GO TO 550
        ALOAD(II)=CLOAD(I,1)+ALOAD(II)
    550 II=II+1
        IF((KCDE.EQ.4).OR.(KCDE.EQ.6)) GO TO 600
        ALOAD(II)=CLOAD(I,2)+ALOAD(II)
        GO TO 600
    560 ALF=CLOAD(I,2)*PI/180.D0
        COSA=DCOS(ALF)
        SINA=DSIN(ALF)
        IIP=II+1
        A1=ALOAD(II)*COSA+ALOAD(IIP)*SINA+CLOAD(I,1)
        A2=ALOAD(IIP)*COSA-ALOAD(II)*SINA
        ALOAD(II)=A1
        ALOAD(IIP)=A2
    CONTINUE
    CALL SOLV(NEQ,NBAND)
    DO 630 I=1,NPBC
        DO 630 J=1,2
            REACT(I,J)=0.0D0
    CALL DISPL
    WRITE(6,1000)
    1000 FORMAT(/, ' DISPLACEMENTS',/,3X, ' JOINT',10X, ' U',22X, ' V',/ )
    DO 700 I=1,NJT
        II=I*2-1
        III=II+1
        700 WRITE(6,2000) I,DISP(II),DISP(III)
        2000 FORMAT(2X,13,2G25.16)
        WRITE(6,3000)
        3000 FORMAT(/, ' REACTIONS',/,3X, ' JOINT',10X, ' RX',22X, ' RY',/ )

```

```

00000860
00000870
00000880
00000890
00000900
00000910
00000920
00000930
00000940
00000950
00000960
00000970
00000980
00000990
00001000
00001010
00001020
00001030
00001040
00001050
00001060
00001070
00001080
00001090
00001100
00001110
00001120
00001130
00001140
00001150
00001160
00001170
00001180
00001190
00001200
00001210
00001220
00001230
00001240
00001250
00001260
00001270
00001280
00001290
00001300
00001310
00001320
00001330

```



```

2500 WRITE(6,2500) REACT(157,1),REACT(157,2)
      FORMAT(5X,2G25.16)
      CALL PUNCH(NJT,DISP)
      CALL STRESS(NPEL)
      GO TO 10
890  END

```

```

SUBROUTINE INPUT
IMPLICIT REAL*8 (A-H,O-Z)
COMMON/INT/NEL,NJT,NMAT,NCLOAD,NPBC,NCON(100,14),NBC(156,2),NSTRES
1  ,NGP,LM(22),LJT(22)
COMMON/FLPL/COORD(156,2),CLOAD(156,2),ELCON(10,2),TITLE(20),T
DATA CHK/,STCP/,
READ (5,1000) TITLE
FORMAT (20A4)
IF (TITLE(1).EQ.CHK) STOP
DO 1050 I=1,156
DO 1050 J=1,2
NBC(I,J)=0
COORD(I,J)=0.0D0
CLOAD(I,J)=0.0D0
DO 1060 I=1,10
DO 1060 J=1,2
ELCON(I,J)=0.0D0
DO 1070 I=1,156
DO 1070 J=1,14
NCON(I,J)=0
1070 WRITE(6,1080) TITLE
1080 FORMAT(1,20A4)
C
C  READ STRUCTURE PARAMETERS
C
READ (5,1100) NEL,NJT,NMAT,NCLOAD,NPBC,NSTRES,NGP,T
FORMAT (7I5,10X,1G24.16)
IF (T.GT.0.0) GO TO 1150
T=1.0D0
1150 WRITE (6,1200) NEL,NJT,NMAT,NCLOAD,NPBC,NSTRES,NGP,T
1200 FORMAT (//,'NUMBER OF ELEMENTS=',I5,/,', TOTAL NUMBER OF JOINTS=',I5,/,',
1  ,I5,/,', NUMBER OF MATERIALS=',I5,/,', NUMBER OF CONCENTRATED LOADS=',I5,/,',
2  ,I5,/,', NUMBER OF JOINTS WITH BOUNDARY CONDITIONS=',I5,/,', THICKNESS OF GAUSS POINTS USED',I5,/,',
3  ,I5,/,', SPECIMEN IS OF THICKNESS',1G24.16)
4  //,
C

```



```

NPBCM=NPBC-1
DO 4000 K=1,NPBC
DO 4000 I=1,NPBC
IF(NBC(I,2).EQ.7) GO TO 3600
GO TO 4000
3600 N1=NBC(I,1)
DO 3800 J=1,NPBCM
NBC(J,1)=NBC(J+1,1)
3800 NBC(J,2)=NBC(J+1,2)
NBC(NPBC,1)=N1
NBC(NPBC,2)=7
4000 CONTINUE
RETURN
END

```

```

SUBROUTINE MERGE(NPEL,NEQ)
IMPLICIT REAL*8 (A-H,O-Z)
DIMENSION SOL/BG(312,64),ALOAD(312),DISP(312),ABGN
COMMON/INT/NEL,NJT,NMAT,NCLOAD,NPBC,NCON(100,14),NBC(156,2),NSTRES
1,NNGP,LM(22),LJT(22)
COMMON/FLPL/CJARD(156,2),CLOAD(156,2),ELCON(10,2),TITLE(20),T
COMMON COORD(12,2),ELAST(3,3),SS(36,24),SN(36,24)
REWIND 10
DO 100 I=1,36
DO 100 J=1,24
SS(I,J)=0.0D0
100 SN(I,J)=0.0D0
DO 200 IK=1,NEL
DO 110 NN=1,NPEL
LJT(NN)=NCON(IK,NN)
110 LM(NN)=(NCON(IK,NN)-1)*2
NS=3*NPEL
DO 120 I1=1,NPEL
DO 120 J1=1,NPEL
I2=LJT(I1)
DO 120 J1=1,2
COORD(I1,J1)=COORD(I2,J1)
120 MATN=NCON(IK,14)
N=NCON(IK,13)*8
YM=ELCON(MATN,1)
PRTIO=ELCON(MATN,2)
CALL ELASM(ELAST,YM,PRTIO,NSTRES,T)
CALL QUAD5(STK,AK,B,NS,N,NGP)
WRITE(10) SS,SN
DO 200 I=1,NPEL
DO 200 J=1,NPEL
DO 200 K=1,2

```


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      II=LM(I)+K
      KK=2*(I-1)+K
      DO 200 L=1,2
      JJ=L4(J)+L-I+1
      IF(JJ.LE.0) GO TO 200
      LL=2*(J-1)+L
      BGK(II,JJ)=BGK(II,JJ)+STK(KK,LL)
200  CONTINUE
      SUM=0.0D0
      DO 300 I=1,NEQ
      SUM=SUM+BGK(I,1)
300  ABGN=SUM*1.0D20
      RETURN
      END

```

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```

C      SUBROUTINE ELASM(D,E,PR,NN,TH)
C      IMPLICIT REAL*8 (A-H,O-Z)
C      THIS SUBROUTINE ASSEMBLYS THE ELASTIC MATRIX FOR PLANE
C      STRAIN (NN=1) OR PLANE STRESS (NN=0)
C      DIMENSION D(3,3)
C      DO 20 I=1,3
C      DO 20 J=1,3
C      D(I,J)=0.0D0
C      IF(NN.EQ.0) GJ TO 10
C      FORM ELASTIC MATRIX FOR PLANE STRAIN
C      ER=TH*E/((1.0D0+PR)*(1.0D0-2.0D0*PR))
C      D(1,1)=ER*(1.0D0-PR)
C      D(1,2)=ER*PR
C      D(2,1)=D(1,2)
C      D(2,2)=D(1,1)
C      D(3,3)=ER*(0.5D0-PR)
C      RETURN
C      FORM ELASTIC MATRIX FOR PLANE STRESS
C      ER=TH*E/((1.0D0-PR*PR)
C      D(1,1)=ER
C      D(1,2)=ER*PR
C      D(2,1)=D(1,2)
C      D(2,2)=D(1,1)
C      D(3,3)=ER*(1.0D0-PR)/2.0D0
C      RETURN
C      END

```


C

CODED BY PROFESSOR GILLES CANTIN AT THE NAVAL POSTGRADUATE SCHOOL

```

COMMON COORD(12,2),ELAST(3,3),SS(36,24),SN(36,24)
DATA X2/0.577350269189626D0,-0.577350269189626D0/
DATA A2/1.0D0,1.0D0/
DATA X3/0.774596669241483D0,0.0D0,-0.774596669241483D0/
DATA A3/0.5555555555555556D0,0.888888888888889D0,0.5555555555555555
156D0/
DATA X4/0.861136311594053D0,0.339981043584856D0,-0.339981043584856
1D0,-0.861136311594053/
DATA A4/0.347854845137454D0,0.652145154862546D0,0.652145154862546D
10.0.347854845137454D0/
DATA X5/0.906179845938664D0,0.538469310105683D0,0.0D0,-0.553846931
10105683D0,-0.906179845938664D0/
DATA A5/0.236926885056189D0,0.478628670499366D0,0.568888888888889
1D0.0.478628670499366D0,0.236926885056189D0/
DATA XYL/1.0D0,-1.0D0,-1.0D0,-1.0D0,1.0D0,1.0D0,-1.0D0,-1.0D0/
DATA XYQ/1.0D0,0.0D0,-1.0D0,-1.0D0,-1.0D0,0.0D0,1.0D0,1.0D0,
11.0D0,1.0D0,0.0D0,-1.0D0,-1.0D0,0.0D0/
DATA XYC/1.0D0,0.333333333333333D0,-0.333333333333333D0,-1.0D0,
1-1.0D0,-1.0D0,-1.0D0,-0.333333333333333D0,0.333333333333333D0,
11.0D0,1.0D0,1.0D0,1.0D0,1.0D0,1.0D0,1.0D0,0.333333333333333D0,
1-0.333333333333333D0,-1.0D0,-1.0D0,-1.0D0,-1.0D0,-1.0D0,-1.0D0,
1333D0,0.333333333333333D0/
IF(NGP.EQ.0) VGP=5
DO 100 I=1,N
DO 100 J=1,N
100 STK(I,J)=0.0D0
      JUMP=NGP-1
DO 140 I=1,NGP
GO TO (132,133,134,135),JUMP
132 XI(I)=X2(I)
      AI(I)=A2(I)
GO TO 140
133 XI(I)=X3(I)
      AI(I)=A3(I)
GO TO 140
134 XI(I)=X4(I)
      AI(I)=A4(I)
GO TO 140
135 XI(I)=X5(I)
      AI(I)=A5(I)
140 CONTINUE
DO 150 I=1,NGP
DO 150 J=1,NGP
150 AIA(I,J)=AI(I)*AI(J)
DO 200 I=1,NGP
X=XI(I)
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```

DO 200 J=1,NGP
Y=XI(J)
CALL FORMK(AK,X,Y,B,N)
DO 200 K=1,N
DO 230 L=1,N
200 STK(K,L)=STK(K,L)+AIA(I,J)*AK(K,L)
DO 300 I=1,N
DO 300 J=1,N
300 STK(I,J)=(STK(I,J)+STK(J,I))*0.5D0
NPT=N/2
IGO=NPT/4
DO 1000 I=1,NPT
GO TO (400,500,600),IGO
400 X=XYL(I,1)
Y=XYL(I,2)
GO TO 700
500 X=XYQ(I,1)
Y=XYQ(I,2)
GO TO 700
600 X=XYC(I,1)
Y=XYC(I,2)
700 CALL FORMK(AK,X,Y,B,N)
DO 800 L=1,3
DO 800 M=1,N
AK(L,M)=0.0D0
800 AK(L,M)=AK(L,M)+ELAST(L,NN)*B(NN,M)
II=3*(I-1)
DO 900 J=1,3
IJ=II+J
DO 900 K=1,N
SS(IJ,K)=AK(J,K)
900 SN(IJ,K)=B(J,K)
1000 CONTINUE
RETURN
END

```

```

SUBROUTINE FORMK(AK,X,Y,B,N)
IMPLICIT REAL*8 (A-H,O-Z)

```

C
C THIS SUBROUTINE FORMS THE STIFFNESS AND THE B MATRIX AS FUNCTIONS OF
C XI AND ETA FOR THE THREE DIFFERENT QUADRILATERAL ELEM. IN THIS FAMILY
C

```

DIMENSION AJ(2,2),AJIN(2,2),DNX(2,12),W1(2,12),B(3,24)

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```

1,B1(24,3),AK(24,24)
COMMON COORD(12,2),ELAST(3,3),SS(36,24),SN(36,24)
C INITIALIZE
DO 10 I=1,3
DO 10 J=1,N
10 B(I,J)=0.0D0
ZERO=0.0D0
ONE=1.0D0
TWO=2.0D0
FOUR=4.0D0
NPT=N/2

C FORM (2XNPT) MATRIX OF DERIV. OF THE INTERP. FUNCT. WRT XI AND ETA
C
C
C IGO=N/8
GO TO (20,30,40),IGO
C
C LINEAR FUNCTIONS
20 W1(1,3)=-(ONE-Y)/FOUR
W1(1,4)=-W1(1,3)
W1(1,1)=(ONE+Y)/FOUR
W1(1,2)=-W1(1,1)
W1(2,3)=-(ONE-X)/FOUR
W1(2,4)=-W1(2,3)
W1(2,1)=-W1(2,4)
W1(2,2)=-W1(2,3)
GO TO 50
30 CONTINUE
C
C QUADRATIC FUNCTIONS
TXPY=TWO*X+Y
TXMY=TWO*X-Y
TYPX=TWO*Y+X
TYMX=TWO*Y-X
CMY=ONE-Y
OPY=ONE+Y
OPX=ONE+X
GMX=ONE-X
W1(1,5)=OMY*TXPY/FOUR
W1(1,6)=-OMY*X
W1(1,7)=OMY*TXMY/FOUR
W1(1,8)=OPY*OMY/TWO
W1(1,1)=OPY*TXPY/FOUR
W1(1,2)=-OPY*X
W1(1,3)=OPY*TXMY/FOUR
W1(1,4)=-OPY*OMY/TWO
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```

W1(2,5)=OMX*TPPX/FOUR
W1(2,6)=-OPX*OMX/TWO
W1(2,7)=OPX*TPMX/FOUR
W1(2,8)=-Y*OPX
W1(2,1)=OPX*TPPX/FOUR
W1(2,2)=OPX*OMX/TWO
W1(2,3)=OMX*TPMX/FOUR
W1(2,4)=-Y*OMX
GO TO 50
CONTINUE

```

40

C
C
C

CUBIC FUNCTIONS

```

OPX=ONE+X
OPY=ONE+Y
OMX=ONE-X
OMY=ONE-Y
TPTMNX=3.0D0+2.0D0*X-9.0D0*X**X
TPTMNX=3.0D0+2.0D0*X-9.0D0*X**X
TPTMNY=3.0D0+2.0D0*Y-9.0D0*Y**Y
TPTMNY=3.0D0+2.0D0*Y-9.0D0*Y**Y
TYPO=3.0D0*Y+ONE
TYMO=3.0D0*Y-ONE
TXPO=3.0D0*X+ONE
TXMO=3.0D0*X-ONE
TMTPEX=10.0D0-27.0D0*X**X+18.0D0*X**X-9.0D0*X**X
TMTMEX=10.0D0-27.0D0*X**X+18.0D0*X**X-9.0D0*X**X
TMNPEY=10.0D0-9.0D0*X**X+18.0D0*X**X-27.0D0*Y**Y
TMNMEY=10.0D0-9.0D0*X**X+18.0D0*X**X-27.0D0*Y**Y
TTT=32.0D0
TTN=32.0D0/9.0D0
W1(1,7)=OMY*TMTPEX/TTT
W1(1,8)=-TPTMNX*OMY/TTN
W1(1,9)=TPTMNX*OMY/TTN
W1(1,10)=-OMY*TMTMEX/TTT
W1(1,11)=-OPY*OMY*TYMO/TTN
W1(1,12)=-OPY*OMY*TYPO/TTN
W1(1,1)=-OPY*TMTMEX/TTT
W1(1,2)=TPTMNX*OPY/TTN
W1(1,3)=-TPTMNX*OPY/TTN
W1(1,4)=-OPY*TMTPEX/TTT
W1(1,5)=-OPY*OMY*TYPO/TTN
W1(1,6)=-OPY*OMY*TYMO/TTN
W1(2,7)=OMX*TMNPEY/TTT
W1(2,8)=OPX*OMX*TXMO/TTN
W1(2,9)=-OPX*OMX*TXPO/TTN
W1(2,10)=OPX*TMNPEY/TTT
W1(2,11)=-TPTMNY*OPX/TTN

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W1(2,12)= TMTMNY*OPX/TTN
W1(2,1) =-OPX*TMNMEY/TTT
W1(2,2) =OPX*OMX*TXPO/TTN
W1(2,3) =-OPX*OMX*TXMO/TTN
W1(2,4) =-OMX*TMNMEY/TTT
W1(2,5) =OMX*TMTMNY/TTN
W1(2,6) =-OMX*TPMTMNY/TTN
50 CONTINUE
C
C FORM JACOBIAN OF THE TRANSFORMATION IN AJ(I,J)
C
DO 100 I=1,2
DO 100 J=1,2
AJ(I,J)=ZERO
DO 100 K=1,NPT
100 AJ(I,J)=AJ(I,J)+W1(I,K)*COORD(K,J)
C
C CALCULATE DETERMINANT OF THE JACOBIAN
C
DTJ=AJ(1,1)*AJ(2,2)-AJ(1,2)*AJ(2,1)
C
C INVERT JACOBIAN IN AJIN
C
AJIN(1,1)=AJ(2,2)/DTJ
AJIN(1,2)=-AJ(1,2)/DTJ
AJIN(2,1)=-AJ(2,1)/DTJ
AJIN(2,2)=AJ(1,1)/DTJ
C
C FORM (2XNPT) MATRIX OF DERIVATIVES OF INTERP. FUNCT WRT X AND Y
C
DO 200 I=1,2
DO 200 J=1,NPT
DNX(I,J)=ZERO
DO 200 K=1,2
200 DNX(I,J)=DNX(I,J)+AJIN(I,K)*W1(K,J)
C
C FORM B(I,J) MATRIX
C
DO 300 I=1,NPT
IOD=2*I-1
IEV=2*I
B(1,IOD)=DNX(1,I)
B(2,IEV)=DNX(2,I)
B(3,IEV)=DNX(1,I)
300 B(3,IOD)=DNX(2,I)
C
C FORM STIFFNESS BY THE CONGRUENT TRANSFORMATION BT*D*B
C

```



```

DO 500 I=1,N
DO 500 J=1,3
B1(I,J)=ZERO
DO 500 K=1,3
500 B1(I,J)=B1(I,J)+B (K,I)*ELAST(K,J)
DO 600 I=1,N
DO 600 J=1,N
AK(I,J)=ZERO
DO 600 K=1,3
600 AK(I,J)=AK(I,J)+B1(I,K)*B(K,J)
DO 700 I=1,N
DO 700 J=1,N
AK(I,J)=((AK(I,J))+AK(J,I))/2.0D0)*DTJ
700 AK(J,I)=AK(I,J)
RETURN
END

```

```

SUBROUTINE CHOL(NEQ,NBAND)
IMPLICIT REAL*8 (A-H,O-Z)
C*****
C THIS SUBROUTINE DOES THE CHOLESKI DECOMPOSITION OF A MATRIX (A)
C WHEN (A) IS SYMMETRIC AND BANDED, (A) = {L}*{L}^T
C THE MATRIX (A) MUST BE FURNISHED IN RECTANGULAR FORM A(NEQ,NBAND)
C WITH ALL ITS DIAGONAL ELEMENTS IN THE FIRST COLUMN
C NEQ IS THE TOTAL NUMBER OF EQUATIONS
C NBAND IS THE HALF-BANDWIDTH OF THE SYSTEM
C CODED BY PROF. CANTIN, DECEMBER 1971
C*****
COMMON/SOL/ A(312,64),B(312),X(312),ABGN
ZRO=0.0D0
ANEQ=NEQ
APZRO=ABGN*1.0D-25/ANEQ
DO 300 I=1,NEQ
DIAG=A(I,1)
IF(DIAG.LE.ZRO) GO TO 400
IF(DIAG.LE.APZRO) WRITE(6,1000) I
1000 FORMAT(/,5X,' SYSTEM IS NEARLY SINGULAR AT EQUATION ',I5//)
DIAG=DSQRT(DIAG)
DO 200 J=1,NBAND
A(I,J)=A(I,J)/DIAG
200 DC 260 J=2,NBAND
L=I+J-1
IF(L.GT.NEQ) GO TO 260
AA=A(I,J)
IF(AA.EQ.ZRO) GO TO 260
DO 250 K=J,NBAND

```



```

250 M=1+K-J
260 A(L,M)=A(L,M)-AA*A(I,K)
300 CONTINUE
300 DNM=0.0D0
350 DO 350 I=1,NEQ
    DNM=DNM+1.0D0/A(I,1)
    ANUM=ABGN*1.0D-20
    CNM=ANUM/DNM
2000 WRITE(6,2000) ANUM, DNM, CNM
    FORMAT(//,5X,' CONDITION NUMBER = (' ,1G24.16,')/(' ,1G24.16,')=' ,
    11G22.16,//)
    RETURN
400 WRITE(6,1100) I,DIAG
1100 FORMAT(//5X,' SYSTEM IS SINGULAR AT EQUATION ',I5,/,5X,' OR INDEFI
1 NITE, DIAG = ',1PG24.16,//)
    STOP
    END

```

```

SUBROUTINE SOLV(NEQ,NBAND)
IMPLICIT REAL*8 (A-H,O-Z)
C*****
C THIS SUBROUTINE PERFORMS A FORWARD SUBSTITUTION FOLLOWED BY A
C BACKWARD SUBSTITUTION IN THE SYSTEM (A)*(X) = (B) WHERE THE MATRIX
C (A) HAS ALREADY BEEN DECOMPOSED BY THE SUBROUTINE CHOL
C (A) IS THE MATRIX OF COEFFICIENTS AFTER A CALL TO S.R.CHOL
C X IS A VECTOR THAT CONTAINS THE ANSWERS UPON EXIT FROM THE S.R.
C NEQ AND NBAND ARE THE SAME AS IN CHOL
C
C AFTER ONE CALL TO S.R.CHOL MANY CALLS TO SOLV CAN BE MADE WITH
C DIFFERENT VALUES FOR THE VECTOR (B), (B) IS DESTROYED EVERY TIME
C CODED BY PROF. CANTIN, DECEMBER 1971
C*****
COMMON/SOL/ A(312,64),B(312),X(312),ABGN
ZRO=0.0D0
DO 200 I=1,NEQ
    IF(B(I).EQ.ZRO) GO TO 200
    B(I)=B(I)/A(I,1)
    DO 100 J=2,NBAND
        L=I+J-1
        IF(L.GT.NEQ) GO TO 100
        AA=A(I,J)
        B(L)=B(L)-AA*B(I)
100 CONTINUE
200 X(NEQ)=B(NEQ)/A(NEQ,1)

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```
DO 400 L=2, NEQ
K=NEQ-L+1
SUM=ZRO
DO 300 J=2, NBAND
M=J+K-1
IF(NEQ.LT.M) GO TO 300
SUM=SUM+A(K,J)*X(M)
CONTINUE
300 X(K)=(B(K)-SUM)/A(K,1)
400 RETURN
END
```

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```
SUBROUTINE STRESS(NPEL)
IMPLICIT REAL*8 (A-H,O-Z)
COMMON/SOL/BGK(312,64),ALOAD(312),DISP(312),ABGN
COMMON/INT/NEL,NJT,NMAT,NCLOAD,NPBC,NCON(100,14),NBC(156,2),NSTRES
1,NGP,LM(22),LJT(22)
COMMON/FLPL/COORD(156,2),CLOAD(156,2),ELCON(10,2),TITLE(20),T
COMMON COORD(12,2),ELAST(3,3),SS(36,24),SN(36,24)
DIMENSION SSJNT(100,7),SNJNT(100,7),SSEL(36),DSPEL(24)
EQUIVALENCE (BGK(1,1),SSJNT(1,1)),(BGK(1,5),SNJNT(1,1)),(BGK(1,9),
1SSEL(1)),(BGK(1,10),SSEL(1)),(BGK(1,11),DSPEL(1))
PI=3.141592653589793D0
REWIND 10
```

DO 50 I=1, NJT

DO 50 J=1, 7

SNJNT(I,J)=0.0D0

SSJNT(I,J)=0.0D0

DO 300 I=1, NEL

READ(10) SS,SN

DO 100 J=1, NPEL

JA=NCON(I,J)*2-1

JJ=J*2-1

DSPEL(JJ)=DISP(JA)

DSPEL(JJ+1)=DISP(JA+1)

NJLM=NPEL*2

NPELT=NPEL*3

DO 200 I1=1, NPELT

SSEL(I1)=0.0D0

SNEL(I1)=0.0D0

DO 200 J1=1, NJLM

SSEL(I1)=SSEL(I1)+SS(I1,J1)*DSPEL(J1)

SNEL(I1)=SNEL(I1)+SN(I1,J1)*DSPEL(J1)

IF(NMAT.EQ.1) GO TO 220

WRITE(6,4000) I

ILO=0

IUP=0


```

DO 210 J=1,NPEL
  IJT=NCON(I,J)
  ILO=IUP+1
  IUP=ILO+2
  WRITE(6,5000) IJT,(SSEL(K),K=ILO,IUP),(SNEL(L),L=ILO,IUP)
210 CONTINUE
5000 FORMAT(5X,I5,6G20.13)
4000 FORMAT(/,2X,' S-T-R-E-S-S-E-S / S-T-R-A-I-N-S FOR ELEMENT',I3
1,17X,' SNY',17X,' SNXY',//)
2,17X,' SSX',17X,' SSY',17X,' SSXY',17X,' SNX',
  ICT=0
DO 300 I2=1,NPEL
  I2A=NCON(I,I2)
DO 250 J2=1,3
  ICT = ICT + 1
  SSJNT(I2A,J2)=SSEL( ICT)+SSJNT(I2A,J2)
  SNJNT(I2A,J2)=SNEL( ICT)+SNJNT(I2A,J2)
  SNJNT(I2A,7)=SNJNT(I2A,7)+1.
DO 500 I3=1,NJT
DO 400 J3=1,3
  SSJNT(I3,J3)=SSJNT(I3,J3)/SNJNT(I3,7)
  SNJNT(I3,J3)=SNJNT(I3,J3)/SNJNT(I3,7)
  GAU=SSJNT(I3,3)
  SXPSY=(SSJNT(I3,1)+SSJNT(I3,2))/2.
  SXMSY=(SSJNT(I3,1)-SSJNT(I3,2))/2.
  AINT=DSQRT(SXMSY*SXMSY+GAU*GAU)
  SSJNT(I3,4)=SXPSY+AINT
  SSJNT(I3,5)=SXPSY-AINT
  SSJNT(I3,6)=(SSJNT(I3,4)-SSJNT(I3,5))/2.000
  IF(DABS(SXMSY).LE.1.0D-6) GO TO 420
  RAN=-GAU/SXMSY
  IF(DABS(RAN).LE.1.0D-2) GO TO 450
  RAN=DATAN(RAN)
  GO TO 450
420 RAN=DSIGN(PI,GAU)/2.000
450 SSJNT(I3,7)=RAN*180.0D0/(2.0D0*PI)
  GAU=SNJNT(I3,3)
  SXPSY=(SNJNT(I3,1)+SNJNT(I3,2))/2.0D0
  SXMSY=(SNJNT(I3,1)-SNJNT(I3,2))/2.0D0
  AINT=DSQRT(SXMSY*SXMSY+GAU*GAU)
  SNJNT(I3,4)=SXPSY+AINT
  SNJNT(I3,5)=SXPSY-AINT
  SNJNT(I3,6)=(SNJNT(I3,4)-SNJNT(I3,5))/2.000
  SNJNT(I3,7)=SSJNT(I3,7)
  WRITE(6,1000)
1000 FORMAT(/,1X,' A-V-E-R-A-G-E S-T-R-E-S-S-E-S AT THE JOINTS',
1//,67X,' PRINCIPLE STRESSES',11X,'MAX. SHEAR',8X,'PRINCIPLE AXIS',
  //,2X,' JOINT',4X,' SSX',12X,' SSY',15X,' SSXY',12X,' SSXM',
1

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1,12X,'SSYM',12X,'SSXYM',15X,'TET(DEG.)',/(/)
DO 600 I=1,NJT
  WRITE(6,2000) I,(SSJNT(I,J),J=1,7)
600  WRITE(6,3000)
3000  FORMAT(//,1X,' A-V-E-R-A-G-E S-T-R-A-I-N-S AT THE JOINTS',
1//,67X,'PRINCIPLE STRAINS',11X,'MAX. SHEAR',8X,'PRINCIPLE AXIS',
1//,2X,'JOINT',4X,'SSX',12X,'SSY',15X,'SSXY',12X,'SSXM',
1,12X,'SSYM',12X,'SSXYM',15X,'TET(DEG.)',/(/)
DO 700 I=1,NJT
  WRITE(6,2000) I,(SNJNT(I,J),J=1,6),SSJNT(I,7)
700  FORMAT(2X,I4,6G17.10,2X,1G19.10)
2000  RETURN
END

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SUBROUTINE BCOND(NBAND)
IMPLICIT REAL*8 (A-H,O-Z)
COMMON/FLPL/CLORD(156,2),CLOAD(156,2),ELCON(10,2),TITLE(20),T
COMMON/SOL/BGK(312,64),ALOAD(312),DISP(312),ABGN
COMMON/INT/NEL,NJT,NMAT,NCLORD,NPBC,NCON(100,14),NBC(156,2),NSTRES
1,NGP,LM(22),LJT(22)
PI=3.141592653589793D0
NEQ=NJT*2
DO 200 I=1,NPBC
  IJT=NBC(I,1)
  IEQ=IJT*2-1
  KODE=NBC(I,2)
  IDIA=0
  GO TO (110,120,130,140,150,160,170),KODE
110  IEQ=IEQ+1
112  BGK(IEQ,1)=ABGN
  IF(IDIA.EQ.1) GO TO 132
  GO TO 200
120  GO TO 112
130  IDIA=1
132  IDIA=0
  GO TO 110
  DXY=CLORD(IJT,2)
  IEQ=IEQ+1
140  ALOAD(IEQ)=ALOAD(IEQ)-BGK(IEQ,1)*DXY
142  ALOAD(IEQ,1)=ABGN
143  BGK(IEQ,1)=2,NBAND
  DO 146 J=2,NBAND
    IROW=IEQ-J+1
    ICOL=IEQ-1+J
    IF(IROW.LT.1) GO TO 144
    ALOAD(IROW)=ALOAD(IROW)-BGK(IROW,J)*DXY
144  IF(ICOL.GT.NEQ) GO TO 146

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00008310

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146 ALOAD(ICOL)=ALOAD(ICOL)-BGK(IEQ,J)*DXY
CONTINUE
IF(IDIA.EQ.1) GO TO 162
GO TO 200
150 DXY=CLOAD(IJT,1)
GO TO 143
160 IDIA=1
DXY=CLOAD(IJT,1)
GO TO 143
162 IDIA=0
GO TO 140
170 ALPHA=CLOAD(IJT,2)*PI/180.0D0
CAL=DCOS(ALPHA)
SAL=DSIN(ALPHA)
IEQP=IEQ+1
DO 180 J=3,NBAND
IROW=IEQ-J+2
JM=J-1
IF(IROW.LT.1) GO TO 175
A1=BGK(IROW,JM)*CAL+BGK(IROW,J)*SAL
A2=BGK(IROW,J)*CAL-BGK(IROW,JM)*SAL
BGK(IROW,JM)=A1
BGK(IROW,J)=A2
175 IF((IEQ+J-1).GT.NEQ) GO TO 180
A3=BGK(IEQ,J)*CAL+BGK(IEQ,JM)*SAL
A4=BGK(IEQ,JM)*CAL-BGK(IEQ,J)*SAL
BGK(IEQ,J)=A3
BGK(IEQ,JM)=A4
180 CONTINUE
A1=BGK(IEQ,1)*CAL*CAL+BGK(IEQ,2)*CAL*SAL*2.0D0+BGK(IEQ,1)*SAL*SAL
A2=BGK(IEQ,2)*((CAL*CAL-SAL*SAL)+(BGK(IEQ,1)-BGK(IEQ,1))*SAL*CAL
BGK(IEQ,1)=A1
BGK(IEQ,2)=A2
BGK(IEQ,1)=ABGN
CONTINUE
RETURN
END

200
SUBROUTINE DISPL
IMPLICIT REAL*8 (A-H,O-Z)
COMMON/FLPL/CJARD(156,2),CLOAD(156,2),ELCON(10,2),TITLE(20),T
COMMON/SOL/BGK(312,64),ALOAD(312),DISP(312),ABGN
COMMON/INT/NEL,NJT,NMAT,NCLoad,NPBC,NCON(100,14),NBC(156,2),NSTRES
1,NNGP,LN(22),LJT(22)
DIMENSION REACT(157,2)
EQUIVALENCE (BGK(1,1),REACT(1,1))
PI=3.141592653589793D0
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DO 200 I=1,NPBC
  IJT=NBC(I,I)
  IEQ=IJT*2-1
  KODE=NBC(I,2)
  IDIA=0
  DPBC=0.0D0
  GO TO (110,120,130,140,150,160,170),KODE
110 IEQ=IEQ+1
  JR=2
112 REACT(I,JR)=-DISP(IEQ)*ABGN
  DISP(IEQ)=DPBC
  IF(IDIA.EQ.1) GO TO 132
  IF(IDIA.EQ.2) GO TO 162
  GO TO 200
120 JR=1
  GO TO 112
130 IDIA=1
  JR=1
  GO TO 112
132 IDIA=0
  GO TO 110
140 DPBC=CLOAD(IJT,2)
  GO TO 110
150 JR=1
  DPBC=CLOAD(IJT,1)
  GO TO 112
160 IDIA=2
  JR=1
  DPBC=CLOAD(IJT,1)
  GO TO 112
162 IDIA=0
  DPBC=CLOAD(IJT,2)
  GO TO 110
170 ALPHA=CLOAD(IJT,2)*PI/180.0D0
  CAL=DCOS(ALPHA)
  SAL=DSIN(ALPHA)
  A1=DISP(IEQ)*CAL
  A2=DISP(IEQ)*SAL
  IEQP=IEQ+1
  A3=-DISP(IEQP)*SAL*ABGN
  A4=DISP(IEQP)*CAL*ABGN
  DISP(IEQP)=A1
  DISP(IEQP)=A2
  REACT(I,1)=-A3
  REACT(I,2)=-A4
  CONTINUE
200 REACT(157,1)=0.0D0
  REACT(157,2)=0.0D0

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DO 300 I=1,NP3C
REACT(157,1)=REACT(157,1)+REACT(I,1)
REACT(157,2)=REACT(157,2)+REACT(I,2)
300 RETURN
END

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**          **          **          **          **          **          **          **          **          **
SUBROUTINE PUNCH
**          **          **          **          **          **          **          **          **          **
SUBROUTINE PUNCH PRODUCES DISPLACEMENT OUTPUT ON PUNCHED CARDS
TO BE USED AS DATA INPUT TO PROGRAM ELASTIC
**          **          **          **          **          **          **          **          **          **
**          **          **          **          **          **          **          **          **          **
IMPLICIT REAL*8 (A-H,O-Z)
COMMON/SOL/BG(312,64),ALOAD(312),DISP(312),ABGN
COMMON/INT/NEL,NJT,NMAT,NCLOAD,NPBC,NCON(100,14),NBC(156,2),NSTRES
1,NGP,LM(22),LJT(22)
DO 20 I=1,NJT
  II=I*2-1
  III=II+1
  WRITE(7,10) I,DISP(II),DISP(III)
  FORMAT(110,2324.16)
10 CONTINUE
20 RETURN
END

```

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BRIAN A. EDWARDS
LIEUTENANT, U. S. NAVY

CHARACTERIZATION OF ELASTIC SOLIDS USING FINITE ELEMENT METHODS

U. S. NAVAL POSTGRADUATE SCHOOL
MONTEREY, CALIFORNIA
DECEMBER 1972

[illegible]


```

ELEMENTAL CONNECTIVITY (1-CARD PER ELEMENT)
COL 1-10 ELEMENT ID, NEID (1110)
11-15 JOINT ID, JOINT I OF ELEMENT
16-20 JOINT ID, JOINT J OF ELEMENT
21-25 JOINT ID, JOINT K OF ELEMENT
26-30 JOINT ID, JOINT L OF ELEMENT

COMMON/INT/NCJN(50,4),L1(6),L2(6),NEL,NJT,NEQ,NEID(50),JID(100),IO
JOINT COORDINATES (1-CARD PER JOINT)
COL 1-10 JOINT ID, JID (1110)
11-34 X-COORDINATE OF JOINT (1G24.16)
35-58 Y-COORDINATE OF JOINT (1G24.16)
JOINT FORCES (1-CARD PER JOINT)
COL 1-10 JOINT ID, JID (1110)
11-34 RESULTANT HORIZONTAL FORCE ACTING ON JOINT (1G24.16)
35-58 RESULTANT VERTICAL FORCE ACTING ON JOINT (1G24.16)
JOINT DISPLACEMENTS (1-CARD PER JOINT)
COL 1-10 JOINT ID, JID (1110)
11-34 HORIZONTAL DISPLACEMENT OF JOINT (1G24.16)
35-58 VERTICAL DISPLACEMENT OF JOINT (1G24.16)

***** CODED BY BRIAN A. EDWARDS, LT., USN (AUG. 72) *****
CALL INPUT
CALL MERGE
CALL BMULT
CALL SINGL (BSYM,6,B,L1,L2)
CALL INVRT
CALL ANSER
GO TO 99
END

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99

```

SUBROUTINE INPUT
IMPLICIT REAL*8(A-H,O-Z)
COMMON/INT/NCJN(50,4),L1(6),L2(6),NEL,NJT,NEQ,NEID(50),JID(100),IO
COMMON/FLP/BK(200,6),BKT(6,200),R(200),COORD(100,2),SS(100,6),SN(
1100,6),SK(8,6),BSYM(6,6),ASYM(6,6),U(130),V(100),BR(6),C(6),A,B,
1 TITLE(20),I
DATA CHK/,STOP,/
READ (5,10) TITLE
FORMAT (20A4)
10 IF (TITLE(1).EQ.CHK) STOP
WRITE (6,11) TITLE
FORMAT (1,20A4,/)
11 READ PARAMETERS
READ (5,15) NEL,NJT,T

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15  FORMAT (2I10,10X,1G24.16)
    IF (T.GT.0.0) GO TO 1150
1150  NEQ=NJT*2
      WRITE (6,16) NEL,NJT,T
16  FORMAT(,0,NUMBER OF ELEMENTS IN STRUCTURE',1I10,
1/,',,NUMBER OF JOINTS IN STRUCTURE',1I10,
1/,',,THICKNESS IS(UNIFORM)',1G24.16)
      READ CONNECTIVITY MATRIX
      DO 20 I=1,NEL
20  READ (5,25) NEID(I), (NCON(I,J),J=1,4)
25  FORMAT (1I10,4I5)
      WRITE (6,30)
30  FORMAT(,J,CONNECTIVITY MATRIX',//,6X,'EL.ID',4X,'I-J-K-L CONNECT
1IVITY,')
      DO 35 I=1,NEL
35  WRITE (6,36) NEID(I), (NCON(I,J),J=1,4)
36  FORMAT (,I,1I10,4I5)
      READ COORDINATES OF EACH JOINT
      WRITE (6,46)
46  FORMAT(,0,2X,'JOINT ID',12X,'X-COORDINATE',12X,'Y-COORDINATE')
      DO 40 M=1,NJT
      DO 40 N=1,2
40  COORD(M,N)=0.0D0
      DO 42 M=1,NJT
      READ (5,45) JID(M), (COORD(M,N),N=1,2)
45  FORMAT (1I10,2G24.16)
42  WRITE (6,44) JID(M), (COORD(M,N),N=1,2)
44  FORMAT (,I,1I10,2G24.16)
      READ FORCES ACTING AT EACH JOINT
      WRITE (6,56)
56  FORMAT(,0,2X,',' ,12X,'HORIZ. FORCE',12X,' VERT. FORCE')
      DO 50 M=1,NEQ
      R(M)=0.0D0
      READ (5,55) (R(M),M=1,NEQ)
55  FORMAT (10X,2G24.16)
52  WRITE (6,54) (R(M),M=1,NEQ)
54  FORMAT (,I,10X,2G24.16)
      READ DISPLACEMENTS OF EACH JOINT
      WRITE (6,66)
66  FORMAT(,0,2X,'JOINT ID',10X,'X-DISPLACEMENT',10X,'Y-DISPLACEMENT',
1)
      DO 60 M=1,NJT
      U(M)=0.0D0
      V(M)=0.0D0
60  FORMAT (,I,10X,2G24.16)
      READ (5,65) JID(M),U(M),V(M)
65  FORMAT (1I10,2G24.16)

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62 WRITE(6,64) JID(M),U(M),V(M)
64 FORMAT(' ',110,2G24.16)
END

SUBROUTINE MERGE
IMPLICIT REAL*8(A-H,O-Z)
COMMON/INT/NCON(50,4),L1(6),L2(6),NEL,NJT,NEQ,NEID(50),JID(100),IO
COMMON/FLP/BK(200,6),BKT(6,200),R(200),COORD(100,2),SS(100,6),SN(
1100,6),SK(8,6),BSYM(6,6),ASYM(6,6),U(100),V(100),BR(6),C(6),A,B,
1ITILE(20),T
**
SUBROUTINE MERGE FORMS THE MASTER MODIFIED STIFFNESS MATRIX
BK---MASTER MODIFIED STIFFNESS MATRIX
SK---ELEMENT MODIFIED STIFFNESS MATRIX
NCON---CONNECTIVITY MATRIX
NEL---NUMBER OF ELEMENTS
NEQ---NUMBER OF EQUATIONS, NJT*2
NJT---NUMBER OF JOINTS IN MASTER ELEMENTS AT A JOINT
U,V---HORIZ. AND VERT. DISPLACEMENTS
T---THICKNESS (ASSUMED UNIFORM)
**
DO 100 I=1,NEQ
DO 100 J=1,6
BK(I,J)=J.OO
DETERMINE EACH ELEMENTS MODIFIED STIFFNESS MATRIX
DO 150 I=1,NJT
DO 150 J=1,6
SS(I,J)=O.OO
SN(I,J)=O.OO
DO 250 IO=1,NEL
CALL STIFF
DO 200 M=1,4
IJ=NCON(IO,M)
II=(M*2)-1
JJ=I+1
DO 175 K=1,6
SS(IJ,K)=SS(IJ,K)+SK(IJ,K)
SN(IJ,K)=SN(IJ,K)+SK(JJ,K)
CONTINUE
DO 200 CONTINUE
DO 250 CONTINUE
DO 300 I=1,NJT
I1=(I*2)-1
I2=I+1
DO 275 K=1,6
BK(I1,K)=SS(I,K)*T

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BK(I2,K)=SN(I,K)*T
275 CONTINUE
300 CONTINUE
RETURN
END

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SUBROUTINE STIFF
IMPLICIT REAL*8(A-H,O-Z)
COMMON/INT/NCON(50,4),L1(6),L2(6),NEL,NJT,NEQ,NEID(50),JID(100),ID
COMMON/FLP/BK(200,6),BKT(6,200),R(200),COORD(100,2),SS(100,6),SN(
100,6),SK(8,6),BSYM(6,6),ASYM(6,6),U(100),V(100),BR(6),C(6),A,B,
111ILE(20),T
DIMENSION SK(10)
*****
SUBROUTINE STIFF LOADS THE ELEMENTAL MODIFIED STIFFNESS MATRIX
SK---ELEMENTIVITY MATRIX
NCON=CONNECTIVITY MATRIX
A,B---HORIZ. AND VERT. DIMENSIONS OF AN ELEMENT
U,V---HORIZ. AND VERT. DISPLACEMENTS AT A JOINT
NEL---NUMBER OF ELEMENTS
NJT---NUMBER OF JOINTS IN MASTER STRUCTURE
*****
DO 10 I=1,10
SSK(I)=0.0D0
DO 20 J=1,8
DO 30 K=1,6
SK(I,J)=0.0D0
CALCULATE THE HORIZONTAL DIMENSION, A
A1=COORD(NCON(I,1),1)
A2=COORD(NCON(I,2),1)
DA=DABS(A1-A2)
A1=COORD(NCON(I,4),1)
A2=COORD(NCON(I,3),1)
DB=DABS(A1-A2)
A=(DA+DB)/2.0D0
CALCULATE THE VERTICAL DIMENSION, B
A1=COORD(NCON(I,4),2)
A2=COORD(NCON(I,1),2)
DA=DABS(A1-A2)
A1=COORD(NCON(I,3),2)
A2=COORD(NCON(I,2),2)
DB=DABS(A1-A2)
B=(DA+DB)/2.0D0
DEFINE SOME COMMON MULTIPLICATIVE CONSTANTS
D1=1.0/2.0D0
D2=1.0/3.0D0
D3=1.0/4.0D0

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D4=A/B
D5=B/A
DEFINE THE TERMS OF AUXILIARY ARRAY SSK
SSK(1)=D1*(U(NCON(IO,1))-U(NCON(IO,3)))
SSK(2)=D2*D5*(V(NCON(IO,1))-V(NCON(IO,2))-D1*V(NCON(IO,3))+D1*V(NCON
1CN(IO,4)))
SSK(3)=D1*(-U(NCON(IO,2))+U(NCON(IO,4)))
SSK(4)=D2*D5*(-D1*V(NCON(IO,1))+D1*V(NCON(IO,2))+V(NCON(IO,3))-V(NCON
1CN(IO,4)))
SSK(5)=D1*(V(NCON(IO,1))-V(NCON(IO,3)))
SSK(6)=D2*D4*(U(NCON(IO,1))+D1*U(NCON(IO,2))-D1*U(NCON(IO,3))-U(NCON
1CN(IO,4)))
SSK(7)=D1*(-V(NCON(IO,2))+V(NCON(IO,4)))
SSK(8)=D2*D4*(D1*U(NCON(IO,1))+U(NCON(IO,2))-U(NCON(IO,3))-D1*U(NCON
1CN(IO,4)))
SSK(9)=D3*(V(NCON(IO,1))-V(NCON(IO,2))-V(NCON(IO,3))+V(NCON(IO,4)))
1)
SSK(10)=D3*(U(NCON(IO,1))+U(NCON(IO,2))-U(NCON(IO,3))-U(NCON(IO,4))
1)
LOAD THE ELEMENTAL MODIFIED STIFFNESS MATRIX
SK(1,1)=D2*D5*(U(NCON(IO,1))-U(NCON(IO,2))-D1*U(NCON(IO,3))+D1*U(NCON
1CN(IO,4)))
SK(1,2)=D3*(V(NCON(IO,1))+V(NCON(IO,2))-V(NCON(IO,3))-V(NCON(IO,4))
1)
SK(3,1)=-SK(1,1)
SK(5,1)=D2*D5*(-D1*U(NCON(IO,1))+D1*U(NCON(IO,2))+U(NCON(IO,3))-U(N
1CN(IO,4)))
SK(7,1)=-SK(5,1)
SK(6,6)=-SK(10)+SSK(4)
SK(3,6)=SSK(9)+SSK(8)
SK(2,6)=SSK(10)+SSK(2)
SK(1,6)=SSK(9)+SSK(6)
SK(8,5)=-SSK(7)+SSK(6))
SK(6,5)=-SSK(5)+SSK(8))
SK(4,5)=SSK(7)+SSK(8)
SK(2,5)=SSK(5)+SSK(6)
SK(4,4)=D2*D4*(D1*V(NCON(IO,1))+V(NCON(IO,2))-V(NCON(IO,3))-D1*V(NCON
1CN(IO,4)))
SK(2,4)=D2*D4*(V(NCON(IO,1))+D1*V(NCON(IO,2))-D1*V(NCON(IO,3))-V(NCON
1CN(IO,4)))
SK(7,3)=-SSK(3)+SSK(4))
SK(5,3)=-SSK(1)+SSK(4)
SK(3,3)=SSK(3)-SSK(2)
SK(1,3)=SSK(1)+SSK(2)
SK(2,2)=D3*(U(NCON(IO,1))-U(NCON(IO,2))-U(NCON(IO,3))+U(NCON(IO,4))
1)
SK(3,2)=-SK(1,2)
SK(4,2)=SK(2,2)

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C

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SK(5,2)=-SK(1,2)
SK(6,2)=-SK(2,2)
SK(7,2)=SK(1,2)
SK(8,2)=SK(6,2)
SK(2,3)=SK(1,1)
SK(4,3)=-SK(1,1)
SK(6,3)=SK(5,1)
SK(8,3)=-SK(4,4)
SK(6,4)=-SK(2,4)
SK(8,4)=SK(2,4)
SK(1,5)=SK(4,4)
SK(3,5)=SK(6,4)
SK(5,5)=SK(8,4)
SK(7,5)=-SK(2,6)
SK(4,6)=-SK(3,6)
SK(5,6)=-SK(1,6)
SK(7,6)=-SK(6,6)
SK(8,6)=-SK(6,6)
RETURN
END

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SUBROUTINE BMULT (A-H,O-Z)
IMPLICIT REAL*8(A-H,O-Z)
COMMON/INT/NCJN(50,4),LI(6),L2(6),NEL,NJT,NEQ,NEID(50),JID(100),ID
*****
SUBROUTINE BMULT PRE-MULTIPLIES THE MASTER STIFFNESS MATRIX AND
THE FORCE VECTOR BY THE TRANSPOSE OF BK
BK---MASTER MODE OF THE STIFFNESS MATRIX
BKT---TRANSPOSE OF THE MASTER STIFFNESS MATRIX
BSYM--SYMMETRIC SQUARE MATRIX,BKT*BK
NEQ---NUMBER OF EQUATIONS,NJT*2
P-----FORCE VECTOR
BR-----COLUMN VECTOR,BKT*P
*****
COMMON/FLP/BK(200,6),BKT(6,200),R(200),COORD(100,2),SS(100,6),SN(
1100,6),SK(8,6),BSYM(6,6),ASYM(6,6),U(100),V(100),BR(6),C(6),A,B,
1 TITLE(20),T
1 FORM THE TRANSPOSE OF THE MASTER MODIFIED STIFFNESS MATRIX
DO 10 I=1,6
DO 10 J=1,NEQ
DO 10 BKT(I,J)=0.0D0
DO 20 I=1,6
DO 20 J=1,NEQ
DO 20 BKT(I,J)=BK(J,I)
1 FORM THE PRODUCT, BSYM=BKT*BK
DO 30 I=1,6
DO 30 J=1,6

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```

      BSYM(I,J)=0.0D0
      DO 30 K=1,NEQ
      BSYM(I,J)=BSYM(I,J)+BKT(I,K)*BK(K,J)
30  FORM THE PRODUCT, BR=BKT*BR
      DO 40 I=1,6
      BR(I)=0.0D0
      DO 40 K=1,NEQ
      BR(I)=BR(I)+BKT(I,K)*R(K)
40  RETURN
      END

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      SUBROUTINE SINGL
      IMPLICIT REAL*8(A-H,O-Z)
      COMMON/INT/NCON(50,4),L1(6),L2(6),NEL,NJT,NEQ,NEID(50),JID(100),IO
      COMMON/FLP/BK(200,6),BKT(6,200),COORD(100,2),SS(100,6),SN(
      1100,6),SK(8,6),BSYM(6,6),ASYM(6,6),U(100),V(100),C(6),A,B,
      1TITLE(20),T
      *****
      SUBROUTINE SINGL CALCULATES THE EIGEN VALUES ASSOCIATED WITH BSYM.
      A ZERO EIGEN VALUE INDICATES BSYM TO BE SINGULAR.
      *****

```

C C C C C

```

      *****
      DO 10 I=1,6
      DO 10 J=1,6
      DO 10 ASYM(I,J)=0.0D0
      DO 15 I=1,6
      DO 15 J=1,6
      15 ASYM(I,J)=BSYM(I,J)
      CALL EIGEN (ASYM,6,0,C,A,6)
      WRITE(6,20)
      20 FORMAT(0,'EIGEN VALUES OF BSYM=BKT*BK')
      DO 30 I=1,6
      30 WRITE(6,35) C(I)
      35 FORMAT(' ',10X,I624.16)
      RETURN
      END

```

SUBROUTINE EIGEN (A,N,NOYES,EIVU,EIVR,NDIM)

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C.....
C SUGGESTED FOR USE AT NPG BY PROF. J. G. CANTIN, CODE 59CI
C ORIGINALLY PROGRAMMED BY W. POOLE, COMPUTER CENTER, U.C. BERKELEY
C
C PURPOSE
C THIS SUBROUTINE CALCULATES IN DOUBLE PRECISION ALL THE EIGENVAL-

```



```

C THE CORRECT SUBSPACE AND GIVE FULL DIGITAL INFORMATION. THIS IS
C PERHAPS THE MOST SATISFACTORY FEATURE OF JACOBI'S METHOD. THE
C OTHER TECHNIQUES ... ARE SUPERIOR ... IN SPEED AND ACCURACY, BUT
C IT IS COMPARATIVELY TIRESOME TO OBTAIN ORTHOGONAL EIGENVECTORS
C CORRESPONDING TO PATHOLOGICALLY CLOSE OR COINCIDENT EIGENVALUES. "
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   SYMMETRIC MATRICES", JOURNAL OF ACM, 1, 1959.
2) WILKINSON, J.H., "THE ALGEBRAIC EIGENVALUE PROBLEM", OXFORD
   U. PRESS, 1965, P.277 FF.
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IF(N-1)20,23,21
20 PRINT 22,N
22 FORMAT(4H N=,I3,68H IS TOO SMALL. LIMIT IS 1. RETURN TO CALLING
   1 ROUTINE FROM EIGEN.)
23 PRINT 24, A(1,1)
24 FORMAT(24H IN EIGEN , MATRIX A = ,E14.6)
21 IF(N-160)1,1,3
3 PRINT 2,N
2 FORMAT(3H N=,I5,1 IS TOO LARGE. LIMIT IS 160. RETURN TO CALLING
   1 ROUTINE FROM EIGEN.)
1 RETURN
1 IF(NYES)99,102,99
99 CONTINUE
DO 101 J=1,N
DO 100 I=1,N
EIVR(I,J)=0.0
EIVR(J,J)=1.0
ATOP=0.
DO 112 J=1,N
DO 111 I=1,J
IF(A(I,J)-A(J,I))90,103,90
CONTINUE
A(I,J)=.5*(A(I,J)+A(J,I))
A(J,I)=A(I,J)
CONTINUE
103 IF(ATOP-DABS(A(I,J)))104,111,111
104 ATOP=DABS(A(I,J))
CONTINUE
111 EIVU(J)=A(J,J)
112 IF(ATOP)109,113
109 PRINT 110
110 FORMAT(26H IN EIGEN , MATRIX A = 0 )

```



```

113 RETURN
    AVGF=DFLOAT(N*(N-1))*0.55
    D=0.0
    DO 114 JJ=2,N
    DO 114 II=2,JJ
    S=A(II-1,JJ)/ATOP
114 D=S*S+D
    DSTOP=(1.D-06)*D
    THRESH=DSQRT(D/AVGF)*ATOP
115 IFLAG=0
    DO 130 JCOL=2,N
    JCOL1=JCOL-1
    DO 130 IROW=1,JCOL1
    AIJ=A(IROW,JCOL)
    IF(DABS(AIJ)-THRESH) 130,130,117
117 AIJ=A(IROW,IROW)
    AIJ=A(JCOL,JCOL)
    S=AIJ-AIJ
    IF(DABS(AIJ)-1.D-09*DABS(S)) 130,130,118
118 IFLAG=1
    IF(1.D-10*DABS(AIJ)-DABS(S)) 116,119,119
119 S=7071067811865475
    C=S
    GO TO 120
116 T=AIJ/S
    S=0.25/DSQRT(0.25+T*T)
    C=DSQRT(0.5+S)
    S=2.*T*S/C
120 DO 121 I=1,IROW
    T=A(I,IROW)
    U=A(I,JCOL)
    A(I,IROW)=C*T-S*U
    A(I,JCOL)=S*T+C*U
121 I2=IROW+2
    IF(I2-JCOL) 127,127,123
127 CONTINUE
    DO 122 I=I2,JCOL
    T=A(I-1,JCOL)
    U=A(IROW,I-1)
    A(I-1,JCOL)=S*U+C*T
    A(IROW,I-1)=C*U-S*T
122 A(JCOL,IROW)=S*AIJ+C*AJJ
123 A(IROW,IROW)=C*AIJ+C*AJJ
    DO 124 J=JCOL,N
    T=A(IROW,J)
    U=A(JCOL,J)
    A(IROW,J)=C*T-S*U
    A(JCOL,J)=S*T+C*U
124

```



```

131 IF(NOVES)131,126,131
    CONTINUE
DO 125 I=1,N
  T=EIVR(I,IROW)
  EIVR(I,IROW)=C*T-EIVR(I,JCOL)*S
125 EIVR(I,JCOL)=S*T+EIVR(I,JCOL)*C
126 CONTINUE
  S=A(I,J)/ATOP
  D=D-S*S
  IF(D-DSTOP)1260,129,129
1260 D=0
  DO 128 JJ=2,N
    DO 128 II=2,JJ
      S=A(II-1,JJ)/ATOP
      D=S*S+D
      DSTOP=(1.D-36)*D
128 THRESH=D*SQRT(D/AVGF)*ATOP
129 CONTINUE
130 IF(IFLAG)115,134,115
134 T=A(1,1)
  A(1,1)=EIVU(1)
  EIVU(1)=T
  DO 132 J=2,N
    T=A(J,J)
    A(J,J)=EIVU(J)
    EIVU(J)=T
    DO 132 I=2,J
      A(I-1,J)=A(J,I-1)
132 RETURN
133 END
.....
SUBROUTINE INVRT
PURPOSE
  INVERT A MATRIX
USAGE
  CALL INVRT(A,N,D,L,M)
DESCRIPTION OF PARAMETERS
  A AND D MUST BE REAL*8
  A - INPJT MATRIX, DESTROYED IN COMPUTATION AND REPLACED BY
    RESULTANT INVERSE.
  N - ORDER OF MATRIX A
  D - RESULTANT DETERMINANT
  L - WORK VECTOR OF LENGTH N

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CCCCCCCCCCCCCCCC


```

CCCCCCCCCCCCCCCCCCCC
M - WORK VECTOR OF LENGTH N
REMARKS
MATRIX A MUST BE A GENERAL MATRIX
SUBROUTINES AND FUNCTION SUBPROGRAMS REQUIRED
NONE
METHOD
THE STANDARD GAUSS-JORDAN METHOD IS USED. THE DETERMINANT
IS ALSO CALCULATED. A DETERMINANT OF ZERO INDICATES THAT
THE MATRIX IS SINGULAR.
DOCUMENTATION: SUBROUTINE INVRT IS A COPY OF SUBROUTINE DMINV,
A CATALOGUED LIBRARY ROUTINE FROM THE W.R.CHURCH
COMPUTER CENTER, NAVPGSCOL, MONTEREY, CALIF. 93940
.....
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CCCCCCCC
SUBROUTINE INVRT(A,N,D,L,M)
DIMENSION A(1),L(1),M(1)
DOUBLE PRECISION A,D,BIGA,HOLD
.....
SEARCH FOR LARGEST ELEMENT
D=1.0D0
NK=-N
DO 80 K=1,N
NK=NK+N
L(K)=K
M(K)=K
KK=NK+K
BIGA=A(KK)
DO 20 J=K,N
IZ=N*(J-1)
DO 20 I=K,N
IJ=IZ+I
10 IF(DABS(BIGA)-DABS(A(IJ))) 15,20,20
15 BIGA=A(IJ)
L(K)=I
M(K)=J
20 CONTINUE
INTERCHANGE ROWS
C
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C
25 J=L(K) 35,35,25
   IF(J-K) 35,35,25
   KI=K-N
   DO 30 I=1,N
   KI=KI+N
   HOLD=-A(KI)
   JI=KI-K+J
   A(KI)=A(JI)
   A(JI)=HOLD
30 INTERCHANGE COLUMNS
C
C
35 I=M(K)
   IF(I-K) 45,45,38
38 JP=N*(I-1)
   DO 40 J=1,N
   JK=NK+J
   JI=JP+J
   HOLD=-A(JK)
   A(JK)=A(JI)
   A(JI)=HOLD
40 DIVIDE COLUMN BY MINUS PIVOT (VALUE OF PIVOT ELEMENT IS
   CONTAINED IN BIGA)
C
C
45 IF(BIGA) 48,46,48
46 D=0.0D0
   RETURN
48 IF(I-K) 50,55,50
50 IK=NK+I
   A(IK)=A(IK)/(-BIGA)
55 CONTINUE
C
C
   REDUCE MATRIX
C
60 DO 65 I=1,N
   IK=NK+I
   HOLD=A(IK)
   IJ=I-N
   DO 65 J=1,N
   IJ=IJ+N
   IF(I-K) 60,65,60
60 IF(J-K) 62,65,62
62 KJ=IJ-I+K
   A(IJ)=HOLD*A(KJ)+A(IJ)
65 CONTINUE
C
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```

C
C
C
      DIVIDE ROW BY PIVOT
      KJ=K-N
      DO 75 J=1,N
      KJ=KJ+N
      IF(J-K) 70,75,70
      70 A(KJ)=A(KJ)/BIGA
      75 CONTINUE

      PRODUCT OF PIVOTS
      D=D*BIGA

      REPLACE PIVOT BY RECIPROCAL

      FINAL ROW AND COLUMN INTERCHANGE
      80 A(KK)=1.0/BIGA
      CONTINUE

      K=N
      100 K=(K-1)
      105 IF(K) 150,150,105
      I=L(K)
      108 IF(I-K) 120,120,108
      JQ=N*(K-1)
      JR=N*(I-1)
      DO 110 J=1,N
      JK=JQ+J
      HOLD=A(JK)
      JI=JR+J
      A(JK)=-A(JI)
      110 A(JI) =HOLD
      120 J=M(K)
      125 IF(J-K) 100,100,125
      KI=K-N
      DO 130 I=1,N
      KI=KI+N
      HOLD=A(KI)
      JI=KI-K+J
      130 A(KI)=-A(JI)
      150 A(JI) =HOLD
      GO TO 100
      RETURN
      END

```



```

SUBROUTINE ANSER
IMPLICIT REAL*8(A-H,O-Z)
COMMON/INT/NCON(50,4),L1(6),L2(6),NEL,NJT,NEQ,NEID(50),JID(100),IQ
COMMON/FLP/BK(200,6),BKT(6,200),R(200),COORD(100,2),SS(100,6),SN(
100,6),SK(8,6),BSYM(6,6),ASYM(6,6),U(100),V(100),C(6),A,B,
ITITLE(20),T
** ** ** ** **
SUBROUTINE ANSER OUTPUTS THE SIX MATERIAL CONSTANTS. THE VECTOR
OF MATERIAL CONSTANTS IS GENERATED BY MULTIPLICATION OF THE VECTOR
BR BY THE INVERSE OF BSYM.
** ** ** ** **
DO 10 J=1,6
C(J)=0.0D0
DO 20 K=1,6
DO 30 L=1,6
C(K)=C(K)+BSYM(K,L)*BR(L)
50 WRITE(6,50)
50 FORMAT(1,'THE VECTOR OF MATERIAL CONSTANTS',
1//,'TRANSPPOSE<C11,C12,C13,C22,C23,C33>=')
DO 60 I=1,6
60 WRITE(6,70) C(I)
70 FORMAT(10,'15X,1G24.16)
RETURN
END

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C C C C C C

REPRESENTATIVE OUTPUT: 15 KIP LOAD, 4-NODES. (DLASTIC)

NUMBER OF ELEMENTS IN STRUCTURE 4
 NUMBER OF JOINTS IN STRUCTURE 9
 THICKNESS IS(UNIFORM)= 0.100000000000000000 00

CONNECTIVITY MATRIX

EL.ID	I-J-K-L	CONNECTIVITY
1	2 3 6 5	
2	5 6 9 8	
3	1 2 5 4	
4	4 5 8 7	

JOINT ID	X-COORDINATE	Y-COORDINATE
1	10.000000000000000	0.0
2	5.000000000000000	0.0
3	0.0	0.0
4	10.000000000000000	-5.000000000000000
5	5.000000000000000	-5.000000000000000
6	0.0	-5.000000000000000
7	10.000000000000000	-10.000000000000000
8	5.000000000000000	-10.000000000000000
9	0.0	-10.000000000000000

	HORIZ. FORCE	VERT. FORCE
	0.9103955176022619	3.924156487558248
	0.0	7.151687024883492
	-0.9103955176022620	3.924156487558250
	0.0	0.0
	0.0	0.0
	0.0	0.0
	0.0	-5.000000000000000
	0.0	-5.000000000000000
	0.0	-5.000000000000000

JOINT ID	X-DISPLACEMENT	Y-DISPLACEMENT
1	0.0	0.0
2	0.0	0.0
3	0.0	0.0
4	-0.4929293242488149D-03	-0.2617793655895807D-02
5	0.2014108968973132D-18	-0.2189597399146653D-02
6	0.4929293242488153D-03	-0.2617793655895808D-02
7	-0.1315568119229649D-02	-0.5164481934212524D-02
8	0.4872528994009446D-18	-0.3827868226091310D-02
9	0.1315568119229649D-02	-0.6164481934212524D-02

EIGEN VALUES OF BSYM=BKT*BK

0.2031967468551705D-08
0.2443927086398269D-06
0.3938446247547611D-07
0.1542630641666066D-06
0.3839002888303476D-06
0.3825140881455582D-07

THE VECTOR OF MATERIAL CONSTANTS

TRANSPOSE<C11,C12,C13,C22,C23,C33>=

32147.99848740955
9226.475565886535
-0.5826520139618010D-11
32147.99848740942
-0.3403485991231466D-11
11460.76146076146

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<p>Finite element methods are applied to the problem of characterizing linear, anisotropic elastic solids. The conventional finite element displacement formulation is used to simulate an elastic material in plane stress. An inverted finite element formulation is then applied, and the characterizing six material constants are calculated as numerical results.</p> <p>A possible test device for the experimental characterization of anisotropic solids is postulated, the precision of displacement measurements to be required for such a device being determined by random perturbation analysis. Numerical constants accurate to within three percent are predicted if a precision of one part in eight hundred (1/800) can be measured. Numerical constants accurate to within one percent are predicted if a precision of one part in eight thousand (1/8000) can be measured in the test device.</p>			

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